Spring Theory as an Approach to the Unification of Fields

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Abstract
The cosmological constant is necessary to be retained in Einstein’s field equations with value depending on the mass of the source. An overview of the spring theory in astrophysics and cosmology is included in this paper. In short range force, the two interacting particles are point-like vertices connected by a bosonic spring. We also suspect that electron may contain negative sterile neutrino. The self energy of a point charge is not infinite so that renormalisation is not necessary.

Keywords
Cosmological Constant, Dark Matter, Classical Electrostatics, Short Range Interaction

1. Introduction
Astrophysical standard model has confirmed to accept the cosmological constant in relating to dark energy—undetectable particles same as dark matter. The matter distribution inside the universe is roughly dark fluid 95 percent and normal matter 5 percent. Different notions such as dark energy, dark matter, aether, pure space and others are of the same entity. They are different manifestations of the same dark fluid aether, just like the extension and compression of a spring [1] [2]. The famous Michelson-Morley experiments provide no proofs to decline the existence of aether; moreover, the basic assumptions of these experiments are wrong [3] [4]. The theory in this paper is in fact a three dimensional treatment of the de-Sitter Schwarzschild solution. Simply speaking, a spring term is added into Newton’s gravitation. With such a model, we can derive easily the Hubble’s law, explain the missing mass in the rotation curve of galaxies, and depict the short range interaction in which the two point-like vertices are connected by a bosonic spring. Spring theory is based on two resources; the theoretical Yang’s pure space and the results of the Pound-Rebka experiments on the photonic frequency changes.
along a vertical path.

1.1. Yang’s Pure Space [5]

\[ R_{\mu\nu;\lambda} = R_{\mu;\nu} \]  

(1)

Or, after contraction of \( \mu \) and \( \nu \)

\[ R_{\lambda;\lambda} = R_{\lambda;\lambda} \]  

(2)

Properties of these equations had been studied by various authors [6] [7] [8]. Pavelle [6] pointed out that Yang’s pure space is non-physical unless the cosmological constant remains in Einstein’s field equations which was later verified by Mielke [8]. We begin from the second Bianchi Identity

\[ R_{id;m} + R_{in;k} + R_{ik;l} = 0. \]  

(3)

For \( n = l \), we have

\[ R_{ik;m} - R_{im;k} = 0 \]  

(4)

Operating \( g^{ik} \) on above, we have

\[ R_{ik} - R_{ik} = 0 \]  

(5)

\[ R_{ik} = \frac{1}{2} R_{ik} \]  

(6)

Comparing with Equation (2), \( R \) is a constant. Now the Einstein’s field equations with the cosmological constant can be written as

\[ R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R + g_{\mu\nu} \Lambda = 0 \]  

(7)

We obtain in the 4-dimensional case

\[ R_{\mu\nu} = g_{\mu\nu} \Lambda \]  

(8)

\[ R = 4 \Lambda \]  

(9)

indicating that Einstein’s case is a special solution of Yang’s pure space where the covariant derivative of the Ricci tensor in Einstein’s case is zero but not in Yang’s case.

Hence, the cosmological constant needs to be retained but to be re-named as spring constant since it behaves like a harmonic oscillator as we can see later. In a 3-dimensional space, a spring term is added into Newton’s law of gravity:

\[ \frac{GM}{r^2} \pm kr = a \left( k \text{ replaces } \Lambda \right) \]  

(10)

where \( k \) is the spring constant of the source while \( \Lambda \) is assigned as the spring constant of the universe which also known as the cosmological constant. Throughout this paper, only 3-dimensional springs are to be considered.

1.2. The Pound-Rebka Experiments

These famous experiments can be found in many textbooks (see Gravitation by Misn-er/Thorne/Wheeler). The main purpose was to measure the frequency changes of pho-
tons under the earth’s gravity. The Jefferson Physical Laboratory at Harvard used a {\textsuperscript{57}}Fe source placed at a height of 22.6 m above the detector.

Data were obtained when the gamma ray dropped onto the detector:

\[
h\Delta \nu = 3.5 \times 10^{-11} \text{ eV}
\]

(11)

\[\Delta = \text{height dropped } 22.6 \text{ m}
\]

(12)

\[E_0 = h\nu_0 = \text{source energy } 14.4 \text{ keV}
\]

(13)

\[a = \frac{c^2\Delta \nu}{v_0\Delta} = 9.67 \text{ m/s}^2
\]

(14)

which is only true at Harvard, or likewise, the state of Massachusetts. In 1965 Pound and Snider refined the apparatus so that the energy shifts on the upward and downward path gave the measured difference of

\[
\frac{\Delta E}{E} \text{ (down)} - \frac{\Delta E}{E} \text{ (up)} = 4.905 \times 10^{-15}
\]

(15)

Since the first term of Equation (15) is known, the second term will immediately yield the deceleration of 9.85 m/s\(^2\). Taking

- the earth’s rotation \(\omega = 7.3 \times 10^{-5}/\text{sec}\)
- the earth’s radius \(r_0 = 6.4 \times 10^6 \text{ m}\)
- the earth’s mass \(= 6 \times 10^{24} \text{ kg}\)
- \(\theta = 42^\circ\)

being the latitude of Massachusetts where Pound and Rebka performed their experiments at Harvard. Upon substituting the acceleration 9.67 m/s\(^2\) and deceleration 9.85 m/s\(^2\) into Equation (10), we obtain the following two equations

\[
\frac{GM}{r_0^2} - \omega^2 r_0 \cos^2 \theta - kr_0 = 9.67 \text{ m/s}^2
\]

(20)

\[
\frac{GM}{r_0^2} - \omega^2 r_0 \cos^2 \theta + kr_0 = 9.85 \text{ m/s}^2
\]

(21)

Thus, the spring constant of the earth \(k = 1.21 \times 10^{-8} /\text{s}^2\) which, as expected, is different from the cosmological constant. The spring term is so-called because it behaves like an harmonic oscillator.

2. An Overview of the Spring Theory in Astronomy and Astrophysics

2.1. Spring of the Earth

From Equation (10), there exists a point or a spherical shell at a distance of \(r = 32000 \text{ km}\) away from earth where the spring and the earth gravity cancel out each other to give \(a = 0\):

\[
\frac{GM}{r^2} - kr = a (r < 32000 \text{ km})
\]

(22)
\[
\frac{GM}{r^2} - kr = 0 \quad (r = 32000 \text{ km}) 
\]
\[
\frac{GM}{r^2} = a \quad (r > 32000 \text{ km}) 
\]

The last equation shows that the spring breaks at the distance of 32,000 km away from us. Equation (10) gives a clear picture of the fifth force different from the Yukawa type [9]-[15]. However, we have pointed out that the Yukawa type of fifth force is non-logical at \( a = 0 \) and cannot predict Equations ((22), (23) and (24)).

2.2. The Spring of the Moon

The almost vacuum lunar surface provides a frictionless condition for a free falling test to verify the existence of the fifth force as well as to obtain the spring of the moon. The total time travelled by a free falling object through a height \( H \) can simply be found as

\[
T = \left( \frac{2H}{g - kH} \right)^{\frac{1}{2}} 
\]

where \( g = 1.627 \text{ m/s}^2 \) is the moon's gravity.

If fifth force does exist, the total time \( T \) must take longer than the classical one without the spring term [16].

2.3. The Spring of the Sun

The Binet Equation (53) yields the solution [17]

\[
u = \frac{1}{r} = \frac{GM}{h^2} + kD \cos \sqrt{1 - \frac{GM^2}{2h^2c^2\phi}} - \frac{k\phi^4}{G^3M^3} 
\]

where \( D \) is a constant. Setting the cosine part to zero, the spring of the sun is

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<th>Table 1. Orbital details of the inner planets.</th>
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<td>( r ) (10^9 m) from sun</td>
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<td>58</td>
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<tr>
<td>Period ( T ) (days)</td>
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<td>( h ) (10^8 m/s)</td>
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\( M = 2 \times 10^{-8} \text{ kg} \), \( h = 2\pi r^2/T \)

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<th>Table 2. Comparison of different ( k ) (1/sec^2) values.</th>
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L. M. Tsang
\[ k = \frac{G^2 M^2}{h^4} \left( \frac{GM}{r^2} - \frac{1}{r} \right) \]  

(27)

Table 1 can be found in many astronomy textbooks. From Table 2, we can see that the average value of \( k \) is higher than those from various authors.

There are two main reasons of difficulty in determining the value of the sun’s \( k \):
- the value of the two terms inside the bracket of Equation (27) is so close to each other.
- planetary interaction has not been taken into account.

However, the spring term of Equation (26) contributes insignificantly in the perihelion shift of planetary motion as well as the bending of light while grazing the sun.

2.4. The Cosmological Constant of the Universe

There are 3 main parameters in any cosmological model, namely the cosmological constant, the Hubble constant and the matter density [22]. In such a large scale structure, 3 dimensional space is sufficient to depict the universe instead of general relativity [23], Milne [24] and McCrea [25] used Newtonian mechanics to derive the cosmological equations while Harrison used the first law of thermodynamics and equations of hydrodynamics [26].

In the beginning, all matters were compressed into a high density lump of universe followed by a release in such a way that all matters were sprung out by the spring(s) as governed by the equation

\[ \Lambda r - \frac{\frac{4}{3} \pi r^3 \rho G}{r^2} = a = v \frac{dv}{dr} \]  

(28)

which is the same approach as Konuschko [27] except the cosmological term was not considered in his paper. Now,

\[ \Lambda \sim 10^{-35}/s^2 \]  

(29)

\[ \text{present density of our universe } \rho \sim 10^{-27} \text{ kg/m}^3 \]  

(30)

\[ G \sim 10^{-11} \text{ m}^3/\text{sec}^2/\text{kg} \]  

(31)

\[ \text{present radius of our universe } R \sim 4 \times 10^{26} \text{ m} \]  

(32)

Equation (28) can be reduced to, upon integration:

\[ v = (\Lambda - G \rho)^{\frac{1}{2}} r = \Lambda^{\frac{1}{2}} r \]  

(33)

which is just the Hubble’s law having the Hubble constant \( H \sim 10^{-17.5}/s \) in agreement with many literatures [28] where they related \( \Lambda \) to \( \Lambda \)CDM model. From the above data, total mass of the universe is \( M \sim 10^{51} \text{ kg} \). Hence, the universe stops to accelerate when

\[ \Lambda r - \frac{GM}{r^2} = 0 \]  

(34)

or \( r \sim 10^{26} \text{ m} \) which is approximately the present radius of our universe. At least it is a comfortable signal to show the tendency of ceasing to accelerate. Interestingly, matter
at the outer rim of the universe exceeds the speed of light, i.e. $\sim 10^{8.5}$ m/s. Superluminal recession of galaxies is acceptable by some cosmologists [29] [30].

2.5. The Missing Mass in the Rotation Curve of Galaxies

It is already known that the cosmological constant is the answer of dark matter [31], or more precisely, variable cosmological constant [32]. This is explicitly referring to the spring constant of the galaxy, but awaiting to be spelt out. Again, in such a large scale of structure, only approximate estimation can be achieved with the following assumptions:

- aberrations in the observed velocity and distance are unavoidable [33] [34].
- the radius $R$ of the cluster and the velocity can be estimated from the rotation curve.
- $1 \text{kpc} = 3 \times 10^{19}$ m.

In the quantum version of the virial theorem, the average value of the operator $T$ in energy eigenstates in one dimension is given by

$$\langle n | T | n \rangle = \frac{1}{2} \langle n | \frac{\partial V}{\partial x} | n \rangle$$

(35)

where $T$ is kinetic energy and $V$ is potential energy. Since the angular velocity of the galaxies is very small: $\omega \sim 10^{-15} / \text{s}$, the virial theorem holds.

We have studied the rotation curves of galaxies in Figure 1 and Figure 2 with the help of the virial theorem.

$$\langle v^2 \rangle = \left( \frac{GM(r)}{r} \right) + \langle k r^2 \rangle$$

(36)

Nearly all these rotation curves yield the same (for detail see [17])

$$M = \frac{v^2 R}{G} \approx 10^{41} \text{ kg}$$

(37)

$$k = 10^{31} / \text{sec}^2$$

(38)

It is clear that each mass has only one unique spring constant assigned to it. Strictly speaking, a flat curve means that the mass is still decreasing depending on $v^2 - kr^2$ in the virial theorem but good approximation of $k$ can be obtained even though $v$ remains constant over several $\text{kpc}'s$. We take two other papers as a comparison. Firstly, we consider Gessner’s paper [37] who used general relativity to investigate 6 NGC’s, found the mass of galaxies $M = 1.8 \times 10^{41}$ to $8 \times 10^{41} \text{ kg}$ and the spring constant $k = 10^{-31}$ to $10^{-33} / \text{sec}^2$. The results matched our work. The second paper belonged to an Indian team led by B. Aryal [38]. They also used general relativity to investigate 15 NGC’s, found $M = 0.13$ to $7.6 \times 10^{40} \text{ kg}$, and the spring constant $k = -2.7$ to $-9 \times 10^{-31} / \text{sec}^2$ which is negative. Based on the information in their paper including the value of $M$, we used the virial theorem to obtain a positive value of $k \sim 10^{-31} / \text{sec}^2$. Once again, in large scale structure, 3 dimension is sufficient to depict astrophysical phenomena.

3. The Electric Field

The electric field energy density $W$ surrounding a charge $q$ is proportional to the square of the field intensity $E$
Since a charge is always accompanied by its electromagnetic mass \( \delta m \), the total mass of a charge particle is \( M = \text{mechanical mass} + \text{electromagnetic mass} \ \delta m \). The above two masses are non-separable from each other. The relationship of the charge and field density is assumed as

\[
2\pi e E a E \rho \nabla \cdot = = (40)
\]

which seems to be reasonable to say the energy density of the source is proportional to the energy density of its surrounding field. Upon integration

\[
1 \int_{0}^{r} = + \int_{0}^{r} \left( 1 + \frac{B}{r} \right)^{-1} dr = \delta mc^2
\]

The total energy

\[
W_{\text{total}} = \frac{q^2}{2} \int_{0}^{r} \frac{1}{r^2} \left( 1 + \frac{B}{r} \right)^{-1} \ dr = \delta mc^2
\]
which is just the Gauss Law except the right hand side of of Equation (42) is not \( m_e \) in the case of electron. Thus \( B \approx \frac{q^2}{2\delta mc^2} \). The field intensity becomes

\[
E = \frac{q}{r^2} \left( 1 + \frac{q^2}{2\delta mc^2 r} \right)^{-1}
\]

(43)

The potential can be written as

\[
\phi = \frac{q}{r} - \frac{q^3}{4\delta mc^2 r^2} + \cdots \quad \text{(long range, }> 10^{-13} \text{ m})
\]

(44)

\[
= \frac{A}{B} \ln \left( 1 + \frac{B}{r} \right) \quad \text{(short range)}
\]

(45)

where \( A \) and \( B \) need to be determined in short range since coupling is involved. Obviously, for \( \delta m = 0 \), \( E = 0 \), showing that the electromagnetic mass always accompanies with the charge. We can use Bohr’s atomic model to find \( \delta m \). The three key equations including the force, energy and the conservation of angular momentum are, respectively

\[
\frac{mv^2}{r} = \frac{e^2}{r^2} - \frac{e^4}{2\delta mc^2 r^3} + \cdots
\]

(46)

\[
E = -13.6 \text{ eV} = \frac{mv^2}{2} - \frac{e^2}{r} + \frac{e^4}{4\delta mc^2 r^2} + \cdots
\]

(47)

and

\[
mvr = h
\]

(48)

Among the 4 pairs solution after solving the above, the most logical pair is \( r = 5.29 \times 10^{-11} \text{ m} \), \( \delta m = -2.2 \times 10^{-31} \text{ kg} \). There were queries about the internal structure of electron in the last century [39]. Bonnor even raised the question “Does electron contain negative mass?” [40]. As already known [41] that the electron mass \( (0.51 \text{ MeV}/c^2) \) is totally electromagnetic but the radius of the electron is not

\[
R \sim \frac{e^2}{m_e c^2} \sim 2.8 \times 10^{-15} \text{ m}
\]

(49)

Instead, through the electron-positron scattering, the upper limit of \( R \sim 2 \times 10^{-18} \text{ m} \) implies that there is a complex internal structure. The Bohr model provides rooms for the electromagnetic mass but not the spring term. Most probably the spring of the proton at the nucleus breaks off before reaching the orbiting electron. It is yet unknown of how much energy is required to bring this negative mass into the physical world. Another explanation of the negative energy particle is to connect it with anti-particles, or most likely the anti-electron neutrinos having positive energy. The expectation value of the additional term in Equation (47) will produce a perturbation term of

\[
\Delta E_{\text{nlm}} = \frac{e^4}{4\delta mc^2} \langle nlm | \frac{1}{r^2} | nlm \rangle
\]

(50)

resulting to a value less than 1 eV: too small to affect the fine structure of hydrogen
spectrum.

4. The Gravitational Field

As both the Coulomb and Newton’s inverse square law are analogous to one another, the gravitational field from Equation (42) becomes

\[
\frac{GM^2}{2} \int_0^\infty \frac{1}{r^2} \left(1 + \frac{B}{r}\right)^2 \, dr = Mc^2
\]

(51)

where \( B = \frac{GM}{2c^2} \). The field intensity Equation (43) in gravitation becomes

\[
\frac{GM}{r^2} - \frac{G^2M^2}{2c^2r^3} = a \quad \text{(long range)}
\]

(52)

Including the spring term, the new Binet equation can be written as

\[
\frac{d^2u}{d\phi^2} + u = \frac{GM}{\hbar^2} - \frac{G^2M^2u}{2c^2\hbar^2} + \frac{k}{\hbar^2u^3}
\]

(53)

To solve for the above Equation (53), we followed the same procedures as in [42] and [43] to get Equation (26). Comparing the tests with general relativity, the spring term contributes insignificantly in the bending of light while grazing the sun whereas the peri-helion shift of a planet gives

\[
\delta \phi = 2\pi \left(1 - \frac{G^2M^2}{4\hbar^2c^2}\right)
\]

(54)

5. Spring in the Short Range Interaction

Figure 3 shows the quarkonium potential energy which can be governed by the following 3 equations:

a) Cornell potential

\[
V(r) = \frac{a}{r} + br
\]

(55)

where \( a = -0.059 \text{ GeV} \cdot \text{fm}, \ b = 1.16 \text{ GeV/fm} \).

Figure 3. Quarkonium potential from fitting the energy levels [44].
b) Natural log potential

\[ V(r) = a \ln(b r) \]  \hspace{1cm} (56)

where \( a = 0.75 \text{ GeV}, \ b = 4.06 \text{ fm} \).

c) Spring theory (Equation (45) + spring term)

\[ V(r) = a \ln\left(1 + \frac{b}{r}\right) + C \]

\hspace{1cm} (57)

where \( a \) and \( b \) can be estimated roughly from the graph. However, the constant \( C \) is in fact the energy of the spring or rather say, the energy of the confined quarkonia. It follows that

\[ m_{q} c^2 = \frac{1}{2} m_{q} k r^2 = C \] \hspace{1cm} (58)

There are many combinations of \( a \) and \( b \) in Equation (57). For instance, for charmonium, \( C = 3.3 \text{ GeV}, \ a = -1.15 \text{ GeV} \) and \( b = 4.36 \text{ fm} \). For bottomonium, \( C = 9.3 \text{ GeV}, \ a = -2.8 \text{ GeV} \) and \( b = 10.9 \text{ fm} \). The above values are not accurate based on trial and error. However, Equation (57) is the general form for short range interactions. The Cornell potential had been applied in the \( s \)-wave with radial quantum number \( n = 1,2,3,4 \) of heavy quarkonia. By solving the Schrödinger equation, numerous energy eigenvalues are listed [45]. In fact, tracing back to 1981 [46], or even earlier, the Cornell potential was recommended as the unified potential for quarkonia, mesons and baryons. We hereby encourage particle physicists to use spring theory.

6. Discussions

Revisiting the equations from (40) to (45), we come to something interesting:

- total field energy of a charge particle with radius \( R \)

\[ \sim \int_{0}^{R} E^2 \left(4\pi r^2\right) dr = \text{constant} + \text{constant} \left(1 + \frac{B}{R}\right)^{-1} \] \hspace{1cm} (59)

- total self energy of a charge with radius \( R \)

\[ \sim \int_{0}^{R} E^2 \left(4\pi r^2\right) dr = \text{constant} \left(1 + \frac{B}{R}\right)^{-1} \] \hspace{1cm} (60)

For \( R = 0 \), none of the above tends to infinity. In a book written by Sapogin [47], it was mentioned that the classical theory of electromagnetism was fundamentally wrong. The electric field at the centre is zero because \( E \) is a vector. Feyman pointed out that Coulomb’s inverse square law fails at very short distance (see Feyman Lectures on Physics volume 2 chapter 5.8). Hence, renormalisation is not necessary. Perhaps short range Maxwell’s equations can be furtherly elaborated towards a new branch of electromagnetism [48].

References


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