

Effect of Metal Oxide Arrester on Chaotic Behavior of Power Transformers

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Abstract

This paper investigates the Effect of Metal Oxide Arrester (MOV) on the Chaotic Behaviour of power transformers considering nonlinear model for core loss of transformer. The paper contains two parts. In part (1): effect of nonlinear core on the onset of chaotic ferroresonance in a power transformer is evaluated. The core loss is modeled by a third order power series in voltage. in part (2): Effect of Metal Oxide Arrester(MOV) on results in part (1) will be studied. The results reveal that the presence of the arrester has a mitigating effect on ferroresonant chaotic overvoltages. resulted bifurcation diagrams and phase plane diagrams have also been delivered.

Keywords: Chaotic Ferroresonance, MOV, Nonlinear Core Loss, Bifurcation Diagram, Phase Plane Diagram

1. Introduction

Ferroresonance is a complex nonlinear electrical phenomenon that can cause dielectric and thermal problems to components power system. Electrical systems exhibiting ferroresonant behaviour are categorized as nonlinear dynamical systems. Therefore conventional linear solutions can not be applied to study ferroresonance. The prediction of ferroresonance is achieved by detailed modelling using a digital computer transient analysis program [1]. Ferroresonance should not be confused with linear resonance that occurs when inductive and capacitive reactances of a circuit are equal. In linear resonance the current and voltage are linearly related and are frequency dependent. In the case of ferroresonance it is characterised by a sudden jump of voltage or current from one stable operating state to another one. The relationship between voltage and current is dependent not only on frequency but also on other factors such as system voltage magnitude, initial magnetic flux condition of transformer iron core, total loss in the ferroresonant circuit and moment of switching [2]. Ferroresonance may be initiated by contingency switching operation, routine switching, or load shedding involving a high voltage transmission line. It can result in Unpredictable over voltages and high currents. The prerequisite for ferro-

sonance is a circuit containing iron core inductance and a capacitance. Such a circuit is characterized by simultaneous existence of several steady-state solutions for a given set of circuit parameters. The abrupt transition or jump from one steady state to another is triggered by a disturbance, switching action or a gradual change in values of a parameter. Typical cases of ferroresonance are reported in Refs. [4-6]. Theory of nonlinear dynamics has been found to provide deeper insight into the phenomenon. Refs. [7-9] are among the early investigations in applying theory of bifurcation and chaos to ferroresonance. The susceptibility of a ferroresonant circuit to a quasiperiodic and frequency locked oscillations are presented. in Ref [10]. The effect of initial conditions is also investigated. Ref. [11] is a milestone contribution highlighting the effect of transformer modeling on the predicted ferroresonance oscillations. Using a linear model. authors of Ref. [12] have indicated the effect of core loss in damping ferroresonance oscillations. The importance of treating core loss as a nonlinear function of voltage is highlighted in Ref. [13]. An algorithm for calculating core loss from no-load characteristics is given in Ref. [14]. The mitigating effect of transformer with linear core loss model connected in parallel to a MOV arrester is illustrated in Ref. [15]. Ferro resonance in capacitive voltage transformer using nonlinear models

along with (Current/Voltage) limiting filters has been discussed in Refs [16-19]. Evaluation of chaos in voltage transformer, effect of resistance key on the chaotic behavior voltage transformer and subharmonics that produced with ferroresonance in this type transformer are studied in Refs. [20-23]. In all of done researches until now, the effect of Metal Oxide Arrester on the chaotic ferroresonance behaviour of power transformer with nonlinear core loss model have not be done. Therefore, present paper addresses the Effect of Metal Oxide Arrester on chaotic behaviour of power transformer considering nonlinear core loss.

2. Circuit Descriptions and Modeling without MOV

The three-phase diagram for the circuit is shown in **Figure 1**.

The 1100 kV transmission line was energized through a bank of three single-phase as reported in autotransformers Ref. [3]. Ferroresonance occurred in phase A when this phase was switched off on the low-voltage side of the autotransformer; phase C was not yet connected to the transformer at that time. The autotransformer is modeled by a T-equivalent circuit with all impedances referred to the high voltage side. The magnetization branch is modeled by a nonlinear inductance in parallel with a nonlinear resistance which represented by nonlinear saturation characteristic ($\lambda - i_{Lm}$) and nonlinear hysteresis and eddy current characteristics ($v_m - i_{Rm}$), respectively. The hysteresis and eddy current characteristics are calculated from the no-load characteristics by applying the algorithm given in Ref. [14]. The iron core saturation characteristic is given by:

$$i_{Lm} = s_1 \lambda + s_2 \lambda^q \quad (1)$$

For three different values of q , Equation (1) is plotted in **Figure 2**. Coefficients s_1 and s_2 are selected as follows:

$$\text{for } q = 11 \quad s_1 = .0067, s_2 = .0001$$

$$\text{for } q = 7 \quad s_1 = .0067, s_2 = .001$$

$$\text{for } q = 5 \quad s_1 = .0071, s_2 = .0034$$

The exponent q depends on the degree of saturation. It has been found that for adequate representation of the saturation characteristics of a power transformer the exponent q may take the values 5, 7, and 11 in Ref. [8].

According to **Figure 2**, for $q = 11$, flux linkage curve via magnetization current in saturation section has less slope than what can be seen for $q = 5 & 7$. Thus, ferroresonance effects appear earlier. Besides, increasing q will increase both number of stable and unstable points. Then

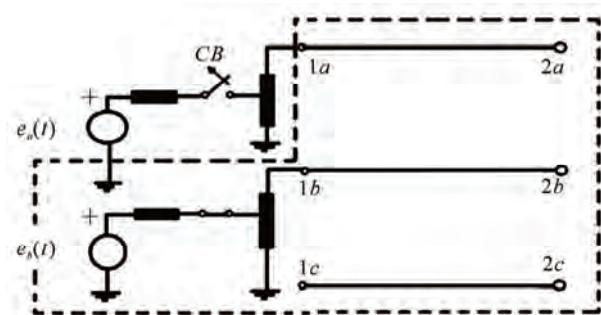


Figure 1. System modeling.

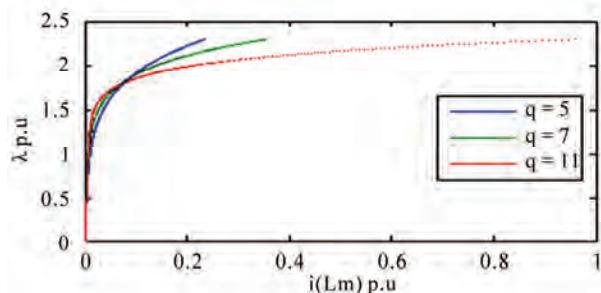


Figure 2. Nonlinear saturation characteristic for three values q .

changing control parameters may expose unstable points like saddle points and chaotic attractors as well.

The core loss is modeled by a switched resistor; which effectively reduced the core loss resistance by a factor of four at the time of onset of ferroresonance. In this paper, the core loss model adopted is described by a third order power series whose coefficients are fitted to match the hysteresis and eddy current nonlinear characteristics given in [3]:

$$i_{Rm} = h_0 + h_1 v_m + h_2 v_m^2 + h_3 v_m^3 \quad (2)$$

Per unit value of (i_{Rm}) for this case given in (3)

$$i_{Rm} = -.000001 + .0047 v_m - .0073 v_m^2 + .0039 v_m^3 \quad (3)$$

Using presented nonlinear functions of core losses and current-flux linkage, plot of i_c (total current) vs. λ in nominal operation for different values of q (5,7,11) are shown in **Figures 4, 5, 6** respectively.

Core current-flux linkage curve includes both eddy current and hysteresis losses. As it can be seen, increasing q will decrease loop width, which consequences reducing losses. Existing even order power loss function causes unsymmetrical loop shape which can be seen in last presented figures.

The circuit in **Figure 1** can be reduced to a simple form by replacing the dotted part with the Thevenin equivalent circuit as shown in **Figure 7**.

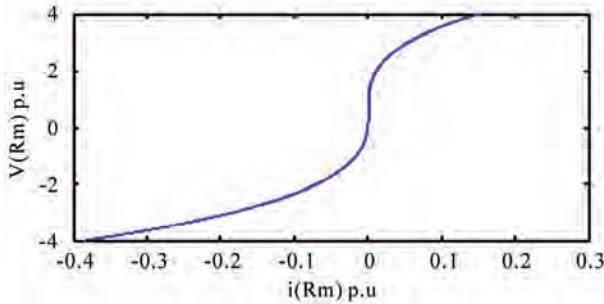
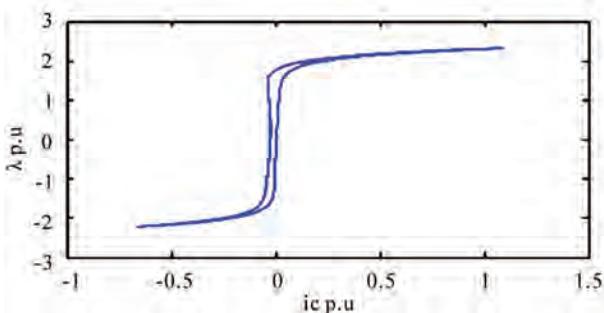
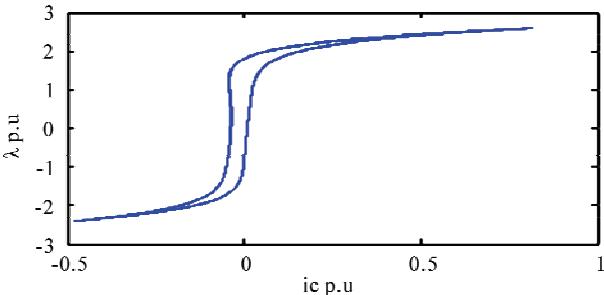


Figure 3. V-I characteristic of nonlinear core loss.

Figure 4. Core current - Flux linkage curve for $q=1$ without MOV and nonnominal operation.Figure 5. Core current - Flux linkage curve for $q=7$ without MOV and nonnominal operation.

Transformer data in [3] Eth and Z_{th} are:

$$E_{th} = 130.1 \text{ kV}; z_{th} = -j1.01243E + 0.5\Omega$$

The resulting circuit to be investigated is shown in **Figure 8** where Z_{th} represents the Thevenin impedance. The behavior of this circuit can be described by the fol-

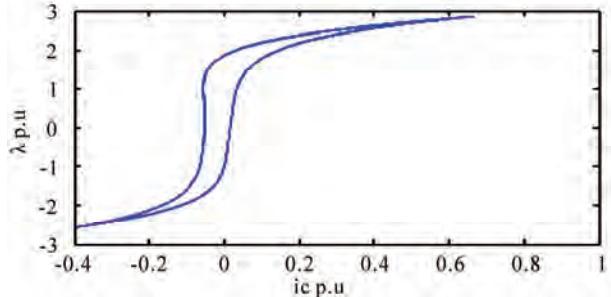
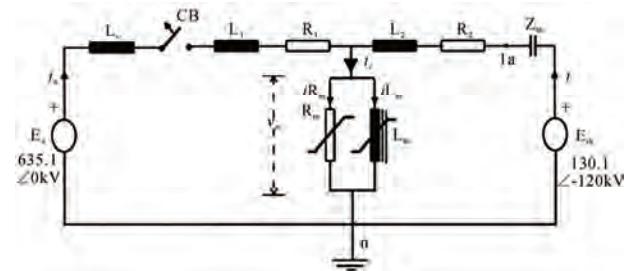
Figure 6. Core current - Flux linkage curve for $q=5$ without MOV and nonnominal operation.

Figure 7. Thevenin circuit of Figure 1.

lowing system of nonlinear differential equations:

The behavior of this circuit can be represented by the following system of nonlinear differential equations:

$$pV_c = \frac{s_1\lambda + s_2\lambda^q + h_0 + h_1P\lambda + h_2(P\lambda)^2 + h_3(P\lambda)^3}{C} \quad (5)$$

With these parameters:

$$\begin{aligned} V_{base} &= 635.1 \text{ kV}, I_{base} = 78.72 \text{ A}, R_{base} = 8067 \Omega, \\ L_{S,p.u.} &= .0188, C_{p.u.} = .07955, \\ R_{p.u.} &= .0014 \quad R_{core} = 556.68 \text{ p.u} \end{aligned}$$

In (4), (5) λ , $P\lambda$ and V_c has been taken as state variables as follows:

$$x_1 = \lambda; x_2 = P\lambda; x_3 = V_c \quad (6)$$

$$Px_1 = x_2 \quad (7)$$

$$Px_3 = \frac{s_1x_1 + s_2x_1^q + h_0 + h_1x_2 + h_2(x_2)^2 + h_3(x_2)^3}{C} \quad (9)$$

$$P^2\lambda = \frac{\left[e_{th}(t) - V_c - P\lambda - R_2 \left(s_1\lambda + s_2\lambda^q + h_0 + h_1(P\lambda) + h_2(P\lambda)^2 + h_3(P\lambda)^3 \right) - L_2 \left(s_1P\lambda + qs_2\lambda^{q-1}(P\lambda) \right) \right]}{\left[L_2 \left(h_1 + 2h_2P\lambda + 3h_3(P\lambda)^2 \right) \right]} \quad (4)$$

$$Px_2 = \frac{\left[e_{th}(t) - x_3 - x_2 - R_2 \left(s_1x_1 + s_2x_1^q + h_0 + h_1(x_2) + h_2(x_2)^2 + h_3(x_2)^3 \right) - L_2 \left(s_1x_2 + qs_2x_1^{q-1}(x_2) \right) \right]}{\left[L_2 \left(h_1 + 2h_2x_2 + 3h_3(x_2)^2 \right) \right]} \quad (8)$$

2.1. Simulation Results and Discussion for Circuit without MOV

Time domain simulations were performed using fourth order Runge-Kutta method and validated against Matlab Simulink results. The initial conditions derived from steady-state solution of Matlab are:

$$x_1 = 0, 0; x_2 = 1.67 \text{ pu}; x_3 = 1.55 \text{ pu}$$

The major analytical tools that in this paper are used to study chaotic ferroresonance are Phase Plane and bifurcation diagram.

The phase plane analysis is a graphical method, in which the time behavior of a system is represented by the movement of state variables of the system in a state space coordination against time. As time evolves, the initial point follows a trajectory. If a trajectory closes on itself, then the system produces a periodic solution. In the chaotic system, the trajectory will never close to itself to shape cycles. A bifurcation diagram is a plot that displays single or multiple solutions (bifurcations) as the value of the control parameter is increased.

Figures 9-11 shows phase plane diagrams of chaotic behaviour for various values of q .

The chaotic behaviour has been intensified with increasing values of q and this behaviour has been shown in these figures.

In **Figures 12-14** core current – flux linkage curves in chaotic ferroresonance mode has been shown. Appearing high frequency oscillations, causes loop width wider than normal operation of power transformer for all q values.

To study voltage effect on system behavior, system has been energized with different supply voltages. **Figures 12-14** show resulted core current-flux linkage curves. All figures are obtained in normal operation of transformer and not in ferroresonance or other abnormal situations. Also **Figures 15-17** show the bifurcation diagram of chaotic behaviours for these three of values of q .

The nonlinear core loss model creates a mitigation on the chaotic ferroresonance behavior in transformer. This

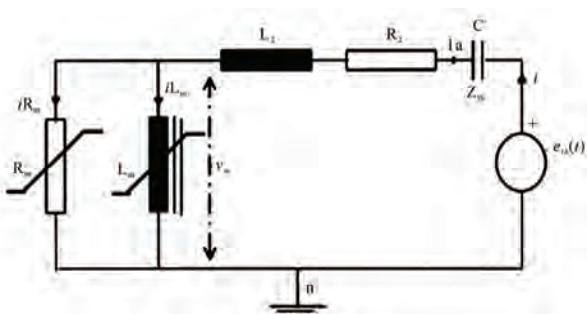


Figure 8. Circuit of ferroresonance investigations without MOV.

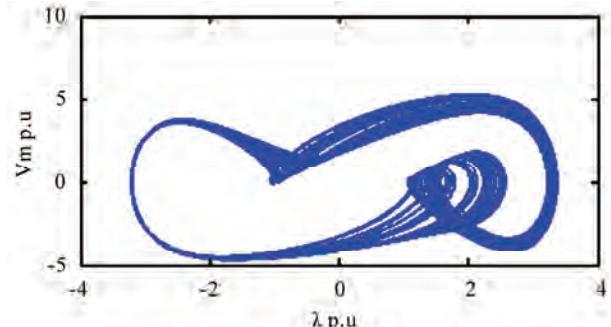


Figure 9. Phase plane diagram for $q = 5$ without MOV.

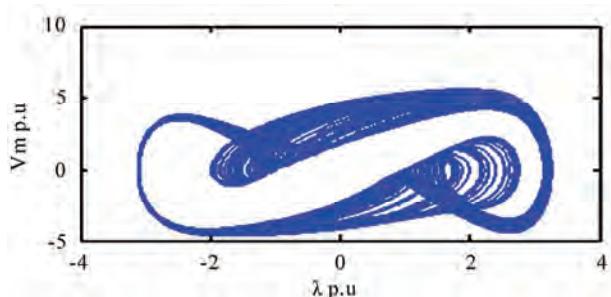


Figure 10. Phase plane diagram for $q = 7$ without MOV.

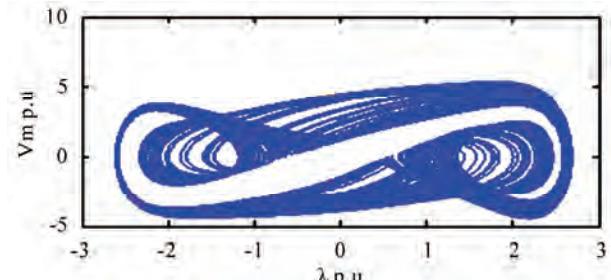


Figure 11. Phase plane diagram for $q = 11$ without MOV.

mitigation is due to delaying terms in nonlinear model. Also presence of nonlinear model for core loss in dynamic equation shows that resulted chaotic ferroresonance behavior of system in comparison with the one with linear core loss model, shows that routing to chaos from period doubling will be smother, more regular and also more distinguishable.

The hatched area under 1.p.u in bifurcation diagrams are due to transmission lines parameters effects. In Ref [8] results has been simplified by elimination of transmission line parameters, which hatched area cannot be seen.

3. Circuit Descriptions and Modeling with MOV

The system considered for analysis consists the MOV

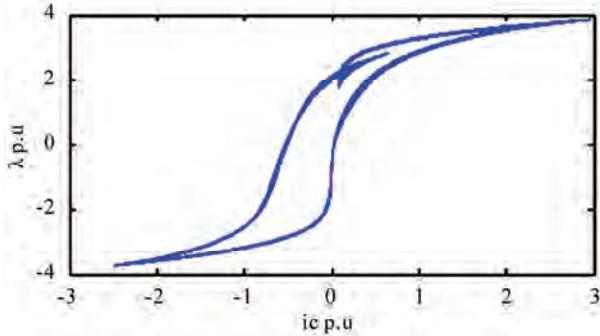


Figure 12. Core current – Flux linkage curve for $q = 5$ without MOV in chaotic condition.

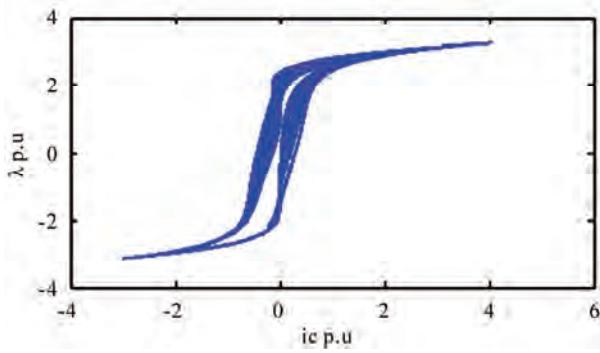


Figure 13. Core current – Flux linkage curve for $q = 7$ without MOV in chaotic condition.

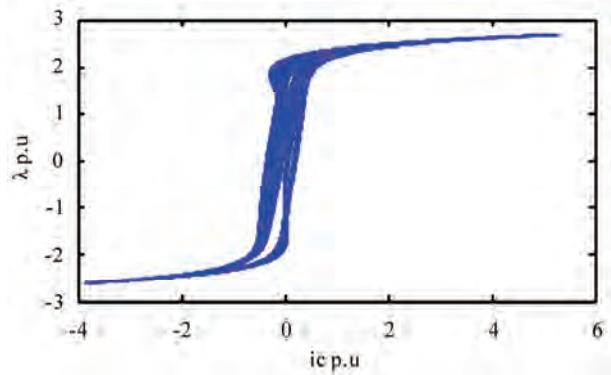


Figure 14. Core current – Flux linkage curve for $q = 11$ without MOV in chaotic condition.

arrester connected across the transformer winding. The related equivalent circuit is as shown in **Figure 18**.

The nonlinear characteristic of the arrester is modelled as [15]:

$$i_{Mov} = \left(\frac{|V_m|}{k} \right)^\alpha \text{sign of } V_m \quad (10)$$

where i_{Mov} , is the current through the arrester in p.u., V_m is the voltage across the transformer winding, k and α are arrester parameters. The differential equation for the circuit in **Figure 18** can be derived as:

$$P^2 \lambda = \frac{\left[e_{th}(t) - V_c - P\lambda - R_2 \left(\left(\frac{d\lambda}{k} \right)^\alpha s_1 \lambda + s_2 \lambda^q + h_0 + h_1(P\lambda) + h_2(P\lambda)^2 + h_3(P\lambda)^3 \right) - L_2(s_1 P\lambda + q s_2 \lambda^{q-1}(P\lambda)) \right]}{\left[L_2 \left(\frac{\alpha}{k} \left(\frac{d\lambda}{k} \right)^{\alpha-1} + h_1 + 2h_2 P\lambda + 3h_3(P\lambda)^2 \right) \right]} \quad (11)$$

$$V_c = \frac{1}{C} \left(\left(\frac{d\lambda}{k} \right)^\alpha + h_0 + h_1(d\lambda) + h_2(d\lambda)^2 + h_3(d\lambda)^3 + a\lambda + b\lambda^q \right) \quad (12)$$

which λ , $P\lambda$ and V_c are assumed as state variables as follows:

$$x_1 = \lambda; x_2 = P\lambda; x_3 = V_c \quad (13)$$

$$Px_1 = x_2 \quad (14)$$

$$dx_2 = - \frac{E - V_c - L_s(ax_2 + qbx_2x_1^{q-1}) - R_s \left(\left(\frac{x_2}{k} \right)^\alpha + h_0 + h_1 + h_2x_2^2 + h_3x_2^3 + ax_1 + bx_1^q \right) - x_2}{L_s \left(\frac{\alpha}{k} \left(\frac{x_2}{k} \right)^{\alpha-1} + h_1 + 2h_2x_2 + 3h_3x_2^2 \right)} \quad (15)$$

$$V_c = \frac{1}{C} \left(\left(\frac{x_2}{k} \right)^\alpha + h_0 + h_1(x_2) + h_2(x_2)^2 + h_3(x_2)^3 + ax_1 + bx_1^q \right) \quad (16)$$

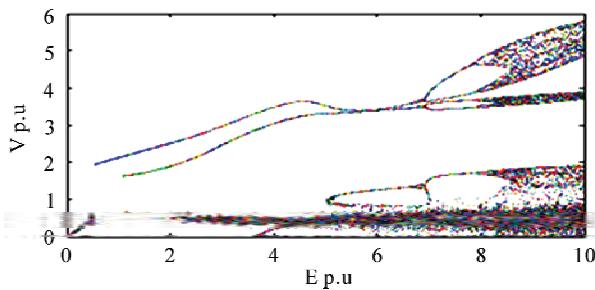


Figure 15. Bifurcation diagram for $q = 5$ without MOV.

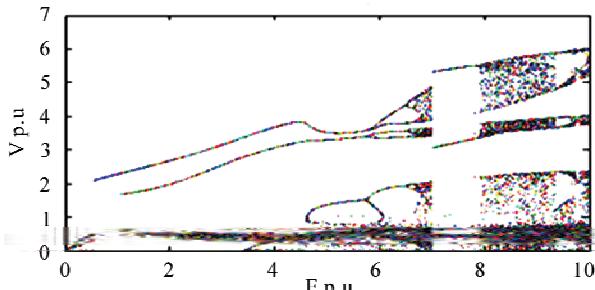


Figure 16. Bifurcation diagram for $q = 7$ without MOV.

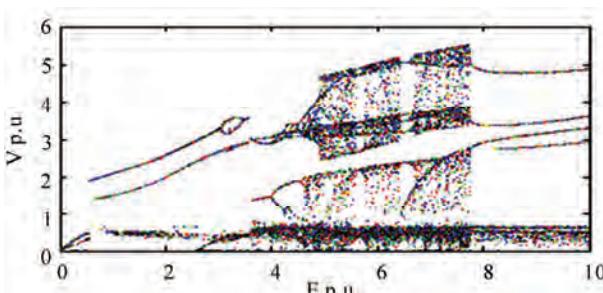


Figure 17. Bifurcation diagram for $q = 11$ without MOV.

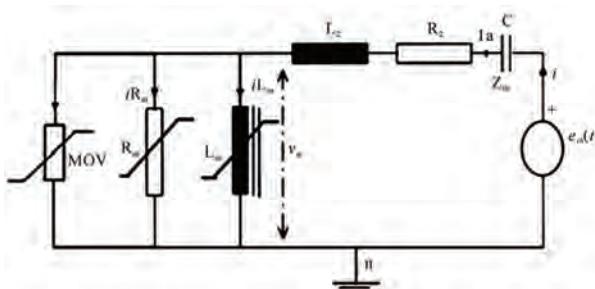


Figure 18. Circuit of ferroresonance investigations with MOV.

The numerical parameters for MOV is:

$$\alpha = 25 \quad k = 2.5101$$

Figures 19, 20, 21 show the time domain simulation phase plane diagram of system states and curve of core current-flux linkage, including arrester for $E = 5$;

The MOV cuts off over voltage so a result of ferro-

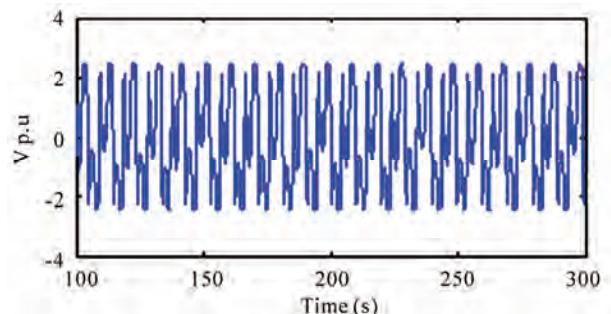


Figure 19. Time domain signal for $E = 5$ and $q = 11$ With MOV.

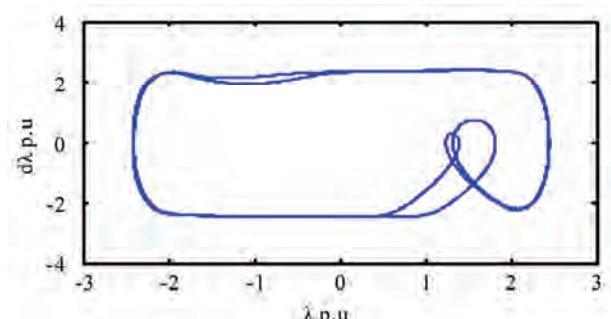


Figure 20. Phase plane diagram for $E = 5$ and $q = 11$ With MOV.

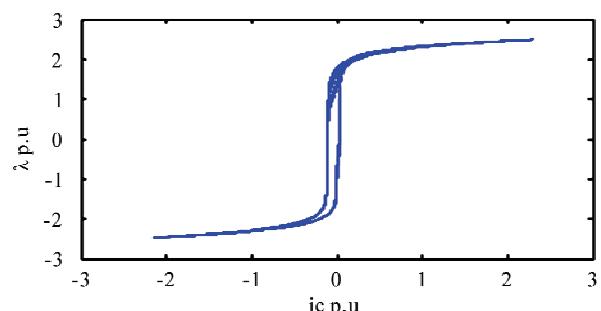


Figure 21. Core current – Flux linkage curve for $q = 11$ while MOV exists in circuit.

sonance chaotic unstable trajectory and eliminates chaos greatly.

Figures 22, 23, and 24 are the bifurcation diagrams by applying MOV surge arrester that show chaotic region mitigates. The tendency for chaos exhibited by the system increases while q increases too

Refer to **Figures 21-24** the MOV in some cases may cause ferroresonance dropout. For $q = 5, 7$ MOV chaotic condition changes to periodic behavior in system, but for $q = 11$ independent chaotic regions which can be created under MOV nominal voltage have survived. Also the pollutions under 1.p.u in bifurcation diagrams due to effect of transmission lines parameters, despite existence of MOV would not be survived.

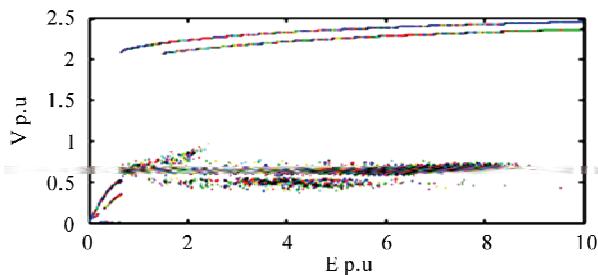


Figure 22. Bifurcation diagram for $q = 5$ with MOV.

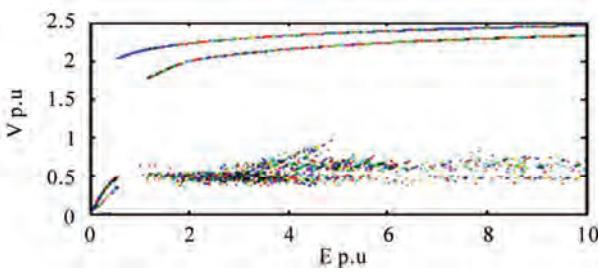


Figure 23. Bifurcation diagram for $q = 7$ with MOV.

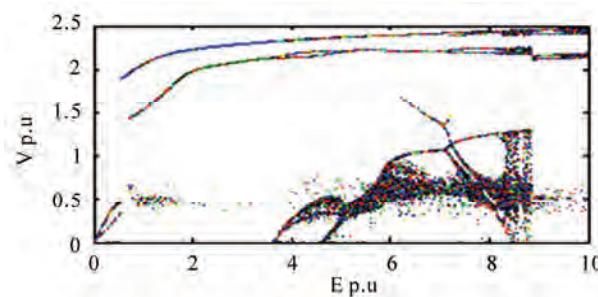


Figure 24. Bifurcation diagram for $q = 11$ with MOV.

4. Conclusions

The dynamic behavior of a transformer is characterized by multiple solutions. The system shows a greater tendency for chaos for saturation characteristics with lower knee points, which corresponds to higher values of exponent q . The presence of the arrester results clamping the Ferroresonant over voltages in studied system. The arrester successfully suppresses or eliminates the chaotic behaviour of proposed nonlinear core loss model. Inclusion of nonlinearity in the core loss reveals that solutions are more realistic and been improved when compared with linear models.

5. References

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Appendix: Nomenclature

- a, b, c index of phase sequence
 h_0, h_1, h_2, h_3 , coefficient for core loss nonlinear function
 n index for the neutral connection
 s_1 coefficient for linear part of magnetizing curve
 s_2 coefficient for nonlinear part of magnetizing curve
 q Index of nonlinearity of the magnetizing curve
 Z_{th} Thevenin's equivalent impedance
 C linear capacitor
 R_m core loss resistance

L nonlinear magnetizing inductance of the transformer

- i instantaneous value of branch current
 v instantaneous value of the voltage across a branch element
 $e_{th}(t)$ instantaneous value of Thevenin voltage source
 e instantaneous value of driving source
 p time derivative operator
 E_{th} R.M.S. value of the Thevenin voltage source
 x state variable
 λ flux linkage in the nonlinear inductance