

Two Very Accurate and Efficient Methods for Solving Time-Dependent Problems

Mohamed El-Gamel, Waleed Adel, M. S. El-Azab

Department of Mathematical Sciences, Faculty of Engineering, Mansoura University, Mansoura, Egypt

Email: gamel_eg@yahoo.com

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Abstract

In this paper, collocation method based on Bernoulli and Galerkin method based on wavelet are proposed for solving nonhomogeneous heat and wave equations. The two methods have the linear systems solved by suitable solvers. Several examples are given to examine the performance of these methods and a comparison is made.

Keywords

Wavelet, Galerkin, Daubechies, Bernoulli, Collocation

1. Introduction

Recently, there has been a great deal of interest in “global” methods (Galerkin and collocation methods) for the numerical solution of two-point boundary value problems. By this, we mean methods which find a solution in the form

$$u_N(x) = \sum_{k=1}^N a_k \psi_k(x) \quad (1)$$

where $\{\psi_k(x)\}_{k=1}^N$ is a basis for an N -dimensional subspace of functions, S . The functions $\psi_k(x)$, $k = 1, 2, \dots, N$, are called test functions and the space S is called the test space.

To simplify the computations, the basis test functions $\{\psi_k(x)\}_{k=1}^N$ are taken to be orthogonal and in many cases they are polynomials, splines, sinc or wavelet. In essence, the Galerkin method is a discretization scheme in which the expansion coefficients $\{a_k\}_{k=1}^N$ are obtained by solving a set of N algebraic equations.

For example, consider the problem $Lu = f$, where L is a self-adjoint operator and f is a known function. The Galerkin method yields the system of equations

$$\sum_{k=1}^N a_k \langle L\psi_k(x), \psi_j(x) \rangle = \langle f(x), \psi_j(x) \rangle, \quad j = 1, \dots, N,$$

which is an algebraic system of equations that can be solved for the unknown coefficients a_k .

In the collocation method no quadrature sums are required. The unknown coefficients a_k in (1) are determined by the linear system

$$Lu_N(x_j) = f(x_j), \quad 0 < x_j < 1, \quad j = 1, 2, \dots, m$$

The aim of this paper is to develop the wavelet and Bernoulli bases for solving PDE of the form

$$\frac{\partial^\alpha u}{\partial t^\alpha} = \frac{\partial^2 u}{\partial x^2} + \eta(x, t), \quad 0 \leq x \leq 1, \quad t > 0, \quad (2)$$

subject to the boundary conditions

$$u(0, t) = h_1(t), \quad u(1, t) = h_2(t), \quad (3)$$

and the initial conditions

$$\begin{aligned} u(x, 0) &= f(x) \\ (\alpha - 1) u_t(x, 0) &= (\alpha - 1) g(x) \end{aligned} \quad (4)$$

where $\alpha \in \{1, 2\}$.

In this paper, there are two different approaches for solving two-point boundary value problems. Galerkin is based on wavelet and collocation is based on using Bernoulli operational matrix to reduce the problem into solving a system of linear algebraic equations. Recently, there has been a lot of research papers dealing with wavelet-Galerkin. Those papers include solution of partial differential equations [1], two-point boundary value problems [2], integro-differential equations [3], second-kind integral equations [5], Fredholm integral equations numerically [4,6], nonhomogeneous time-dependent problems [7], singularly Perturbed convection-dominated diffusion equation [8], telegraph equations [9], eigenvalue problem of a compact integral operator [10], fourth-order multi-dimensional elliptic partial differential equations [11], fourth order linear and nonlinear differential equations [12], stochastic fractional differential equations [13], Schrodinger equations with general nonlinearity [14], generalized wavelet-Galerkin method [15]. El-Gamel *et al.* [16], have compared the wavelet-Galerkin and sinc-Galerkin techniques in solving nonhomogeneous heat equations. Moreover, El-Gamel has compared the wavelet-Galerkin and Adomian decomposition methods of boundary-value problems [17].

Bernoulli matrix method has been used to find the approximate solutions of two-dimensional hyperbolic telegraph equations [18], linear partial differential equations [19], pantograph equation [20], nonlinear fuzzy Hammerstein-Volterra delay integral equations [21], fractional Fredholm-Volterra integro-differential equations [22], the Blasius and MHD Falkner-Skan boundary-layer equations [23], linear multidimensional diffusion and wave equations [24], optimal control problems [25] and Fuzzy integral equations [26]. Recently, El-Gamel and Adel [27] proposed a new approach to solving higher-order boundary value problems via Euler matrix method.

The rest of this article is organized as follows. In Section 2, we describe the wavelet-Galerkin method. In Section 3, Bernoulli-collocation method is introduced. Section 4, gives specific three examples to test the two proposed methods and compare the results. Closing with conclusion In Section 5.

2. Wavelet Bases

2.1. Governing Equation

Setting

$$\frac{\partial^\alpha u}{\partial t^\alpha} \approx \frac{u^{i+1} - \alpha u^i + (\alpha - 1)u^{i-1}}{(\Delta t)^\alpha}$$

Equation (2) may be approximated by

$$\left[\frac{u^{i+1} - \alpha u^i + (\alpha - 1)u^{i-1}}{(\Delta t)^\alpha} \right] = \frac{d^2 u^{i+1}}{dx^2} + \eta(x, t^{i+1}) \tag{5}$$

where $t^i = i\Delta t, i = 0, 1, 2, \dots$. Then, rewriting Equation (5)

$$\frac{d^2 u^{i+1}}{dx^2} - \left[\frac{1}{(\Delta t)^\alpha} \right] u^{i+1} = M(x, t^{i+1}) \tag{6}$$

where

$$M(x, t^{i+1}) = -\eta(x, t^{i+1}) - \left[\frac{\alpha}{(\Delta t)^\alpha} \right] u^i + \left[\frac{(\alpha - 1)}{(\Delta t)^\alpha} \right] u^{i-1}.$$

We assume that

$$\eta(x, t) = h(t)z(x) = \sum_{i=0}^m c_i x^i h(t).$$

2.2. Daubechies Wavelet Bases

More detailed discussions about Daubechies wavelets can be found in [28–31].

2.3. Wavelet-Galerkin Method

Let the solution $u_J(x)$ at the $(i + 1)^{th}$ time level of the problem be approximated by

$$u_J(x) = 2^{J/2} \sum_{k=2-D}^{2^J-1} a_k \phi(2^J x - k), \quad k \in Z, \tag{7}$$

By substituting the solution $u_J(x)$ in Equation (6), yields

$$2^{J/2} \sum_{k=2-D}^{2^J-1} a_k \frac{d^2}{dx^2} [\phi(2^J x - k)] - 2^{J/2} \left[\frac{1}{(\Delta t)^\alpha} \right] \sum_{k=2-D}^{2^J-1} a_k \phi(2^J x - k) = M(x, t^{i+1}). \tag{8}$$

We use the inner product of both sides of Equation (8) with $2^{J/2} \phi(2^J x - l)$ leads the following equation

$$\sum_{k=2-D}^{2^J-1} b_{kl}^J a_k - \left[\frac{1}{(\Delta t)^\alpha} \right] \sum_{k=2-D}^{2^J-1} c_{kl}^J a_k = d_{ml}^J, \quad l = 2-D, 3-D, \dots, 2^J-1. \tag{9}$$

where

$$\begin{aligned} b_{kl}^J &= 2^J \int_0^1 \phi(2^J x - k)\phi(2^J x - l)dx \\ &= \int_{-l}^{2^J-l} \phi(y - (k - l))\phi(y)dy \\ &= \Gamma_{k-l}^0(2^J - l) - \Gamma_{k-l}^0(-l), \end{aligned} \tag{10}$$

$$\begin{aligned}
c_{kl}^J &= 2^{3J} \int_0^1 \phi''(2^J x - k) \phi(2^J x - l) dx \\
&= 2^{2J} \int_{-l}^{2^J - l} \phi''(y - (k - l)) \phi(y) dy \\
&= 2^{2J} [\Gamma_{k-l}^2(2^J - l) - \Gamma_{k-l}^2(-l)],
\end{aligned} \tag{11}$$

and

$$\begin{aligned}
d_{ml}^J &= \int_0^1 \left[-\eta(x, t^{i+1}) - \left[\frac{\alpha}{(\Delta t)^\alpha} \right] u^i + \left[\frac{(\alpha - 1)}{(\Delta t)^\alpha} \right] u^{i-1} \right] \phi_{Jl}(x) dx \\
&= \int_0^1 \left[-h(t^{i+1}) \sum_{i=0}^m c_i x^i - \left[\frac{\alpha}{(\Delta t)^\alpha} \right] u^i \right. \\
&\quad \left. + \left[\frac{(\alpha - 1)}{(\Delta t)^\alpha} \right] u^{i-1} \right] \phi_{Jl}(x) dx \\
&= -h(t^{i+1}) \sum_{i=0}^m \int_0^1 c_i x^i \phi_{Jl}(x) dx - \int_0^1 \left[\left[\frac{\alpha}{(\Delta t)^\alpha} \right] u^i \right. \\
&\quad \left. + \left[\frac{(\alpha - 1)}{(\Delta t)^\alpha} \right] u^{i-1} \right] \phi_{Jl}(x) dx \\
&= -h(t^{i+1}) \sum_{i=0}^m \frac{c_i}{2^{(i+\frac{1}{2})J}} M_l^i(2^J) - \frac{c_{00}}{2^{(\frac{1}{2})J}} M_l^0(2^J).
\end{aligned} \tag{12}$$

The algorithm for calculating Γ_{k-l}^0 , Γ_{k-l}^2 , and M_l^m has been described in [32].

The matrix-vector form of the Equation (9) is

$$\left[W - \left[\frac{1}{(\Delta t)^\alpha} \right] R \right] A = S, \tag{13}$$

where

$$\begin{aligned}
W &= [b_{kl}^J]_{2-D \leq k, l \leq 2^J - 1}, & R &= [c_{kl}^J]_{2-D \leq k, l \leq 2^J - 1}, \\
S &= [d_{ml}^J]_{2-D \leq l \leq 2^J - 1}, & A &= [a_{2-D}, a_{3-D}, \dots, a_{2^J - 1}]^t,
\end{aligned}$$

where t denotes the matrix transpose. Now we have a linear algebraic system that can be solved by the Q-R method.

3. Bernoulli Bases

More detailed discussions about Bernoulli operational matrix can be found in [18, 22, 23, 25].

Bernoulli-Collocation Method

Let the solution of (6) is

$$u_N(x, t^{i+1}) = u_N(x)^{i+1} \simeq \sum_{n=0}^N a_n B_n(x) = B(x) A \tag{14}$$

where

$$\begin{aligned}
A^t &= [a_0, a_1, \dots, a_N] \\
B(x) &= [B_0(x), B_1(x), \dots, B_N(x)]
\end{aligned}$$

then the matrix form of the second derivative is

$$\frac{d^2}{dx^2} u_N(x, t^{i+1}) = \frac{d^2}{dx^2} u_N(x)^{i+1} = B^{(2)}(x) A = B(x) (M^t)^2 A. \quad (15)$$

where M is $(N + 1) \times (N + 1)$ Bernoulli operational matrices of differentiation described by

$$M = \begin{bmatrix} 0 & 0 & \dots & 0 \\ 1 & 0 & \dots & 0 \\ 0 & 2 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & \dots & & N & 0 \end{bmatrix} \quad \text{and} \quad B(x)^t = \begin{bmatrix} B_0(x) \\ B_1(x) \\ \vdots \\ B_{N-1}(x) \\ B_N(x) \end{bmatrix}$$

By replacing each term in Equation (6) with the approximation defined in Equation (14) and (15) and collocate them at $x = x_k$ defined as the equal collocation points where

$$x_k = \frac{k}{N}, \quad k = 0, 1, 2, \dots, N.$$

We reach the following theorem

Theorem 1. *If the assumed approximate solution of the boundary-value problem (6) is (14), then the discrete Bernoulli system is given by*

$$\begin{aligned} & \left[\sum_{n=0}^N B_n''(x_k) - \frac{1}{(\Delta t)^\alpha} \sum_{n=0}^N B_n(x_k) \right] a_n \\ & = -\eta(x_i, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_k, t^i) + \frac{(\alpha - 1)}{(\Delta t)^\alpha} u(x_k, t^{i-1}). \quad (16) \end{aligned}$$

Proof. By replacing each term of Equation (6) with the corresponding approximation represented in Equation (14) and (15) and collocate them with $x = x_k$ collocation points. □

The fundamental matrix for the above system is

$$\Phi A = F$$

where

$$\Phi = B (M^t)^2 - JB,$$

and

$$B = \begin{pmatrix} B_0(x_0) & B_1(x_0) & B_2(x_0) & \dots & B_N(x_0) \\ B_0(x_1) & B_1(x_1) & B_2(x_1) & \dots & B_N(x_1) \\ B_0(x_2) & B_1(x_2) & B_2(x_2) & \dots & B_N(x_2) \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ B_0(x_N) & B_1(x_N) & B_2(x_N) & \dots & B_N(x_N) \end{pmatrix},$$

$$F = \begin{pmatrix} -\eta(x_0, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_0, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_0, t^{i-1}) \\ -\eta(x_1, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_1, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_1, t^{i-1}) \\ \vdots \\ -\eta(x_N, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_N, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_N, t^{i-1}) \end{pmatrix}$$

$$J = \begin{pmatrix} \frac{1}{(\Delta t)^\alpha} & 0 & 0 & \dots & 0 \\ 0 & \frac{1}{(\Delta t)^\alpha} & 0 & \dots & 0 \\ 0 & 0 & \frac{1}{(\Delta t)^\alpha} & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & \frac{1}{(\Delta t)^\alpha} \end{pmatrix}.$$

Then, substituting the approximation from (14) in the boundary conditions in Equation (3) yields

$$\sum_{n=0}^N a_n B_n(0) = h_1(t^{i+1}) \quad (17)$$

and

$$\sum_{n=0}^N a_n B_n(1) = h_2(t^{i+1}). \quad (18)$$

Replacing the first and the last row in the augmented matrix $[\Phi; F]$ with the boundary conditions from Equations (17) and (18) will lead to the new augmented matrix $[\Theta; \tilde{F}]$ which is an $N + 1$ linear equations in $N + 1$ unknowns defined as

$$\Theta A = \tilde{F}$$

where

$$\Theta = \begin{pmatrix} B_0(0) & B_1(0) & B_2(0) & \dots & B_N(0) \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \dots & \dots & \Phi & \dots & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ B_0(1) & B_1(1) & B_2(1) & \dots & B_N(1) \end{pmatrix},$$

$$\tilde{F} = \begin{pmatrix} h_1(t^{i+1}) \\ -\eta(x_0, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_0, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_0, t^{i-1}) \\ -\eta(x_1, t^{i+1}) - 1 \frac{\alpha}{(\Delta t)^\alpha} u(x_1, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_1, t^{i-1}) \\ \vdots \\ -\eta(x_N, t^{i+1}) - \frac{\alpha}{(\Delta t)^\alpha} u(x_N, t^i) + \frac{(\alpha-1)}{(\Delta t)^\alpha} u(x_N, t^{i-1}) \\ h_2(t^{i+1}). \end{pmatrix}$$

This system is solved using the Q-R method for finding the unknown coefficients $\{a_n\}_{n=0}^N$, but with known approximate $u(x_k, t^{i-1})$ and then $u(x_k, t^i)$ can be evaluated, respectively.

4. Numerical Examples

Three examples are considered to demonstrate the efficiency and accuracy of the proposed methods in homogeneous and nonhomogeneous boundary conditions. Daubechies 6 wavelet, $D = 6$, is used and each example was run for $J = 9$ and we take $\Delta t = 0.001$. The maximum absolute error is taken as

$$\|E_{WG}\| = |u_{\text{exact}} - u_{\text{wavelet-Galerkin}}|,$$

and

$$\|E_{BC}\| = |u_{\text{exact}} - u_{\text{Bernoulli-collocation}}|.$$

Example 1 [17] Consider the following problem

$$u_t = u_{xx} \quad 0 \leq x \leq 1, t > 0, \quad (19)$$

subject to the boundary conditions

$$u(0, t) = \exp(t), \quad u(1, t) = \exp(t + 1),$$

and the initial condition

$$u(x, 0) = \exp(x),$$

whose exact solution is

$$u(x, t) = \exp(t + x).$$

Table 1 shows the comparison between the absolute errors $\|\mathbf{E}_{BC}\|$ and $\|\mathbf{E}_{WG}\|$.

Table 1. Maximum absolute errors for Example 1.

(x, t)	$\ \mathbf{E}_{BC}\ , N = 10$	$\ \mathbf{E}_{WG}\ , J = 9.$
(0.25, 0.1)	4.631913E-06	8.5436E-05
(0.5, 0.1)	6.838514E-06	9.3299E-05
(0.75, 0.1)	6.751201E-06	1.7983E-05
(0.25, 0.5)	5.340952E-05	5.4634E-04
(0.5, 0.5)	7.813925E-05	7.6003E-04
(0.75, 0.5)	6.350286E-05	5.7020E-04
(0.25, 1.0)	1.769145E-05	1.3965E-03
(0.5, 1.0)	2.588038E-05	1.7219E-03

Example 2 [17] Consider the following problem

$$u_t = u_{xx} \quad 0 \leq x \leq 1, \quad t > 0, \quad (20)$$

subject to the boundary conditions

$$u(0, t) = 0, \quad u(1, t) = \sin 1 e^{-t},$$

and the initial conditions

$$u(x, 0) = \sin x$$

whose exact solution is

$$u(x, t) = \sin x e^{-t}.$$

The computational results are summarized in **Table 2**.

Table 2. Comparison between the maximum absolute errors for Example 2.

(x, t)	$\ \mathbf{E}_{BC}\ , N = 10$	$\ \mathbf{E}_{WG}\ , J = 9.$
(0.25, 0.1)	1.00112E-06	3.2750E-05
(0.5, 0.1)	1.70939E-06	5.5452E-05
(0.75, 0.1)	1.58965E-06	4.3456E-05
(0.25, 0.5)	6.18197E-06	1.9567E-05
(0.5, 0.5)	9.80818E-06	5.4572E-04
(0.75, 0.5)	8.19902E-06	7.7654E-04
(0.25, 1.0)	7.60298E-06	1.6562E-04
(0.5, 1.0)	1.20468E-05	3.7982E-03
(0.75, 1.0)	1.00501E-05	4.4324E-03

Example 3 [17] Consider the following problem

$$u_{tt} = u_{xx} + 2 \exp(-\pi t) \sin(\pi x), \quad 0 \leq x \leq 1, t > 0, \quad (21)$$

subject to the boundary conditions

$$u(0, t) = 0, \quad u(1, t) = 0,$$

and the initial conditions

$$\begin{aligned} u(x, 0) &= \sin(\pi x) \\ u_t(x, 0) &= -\pi \sin(\pi x) \end{aligned}$$

whose exact solution is

$$u(x, t) = \exp(-\pi t) \sin(\pi x).$$

The computational results are summarized in Table 3.

Table 3. Comparison between the maximum absolute errors for Example 3.

(x, t)	$\ \mathbf{E}_{BC}\ , N = 10$	$\ \mathbf{E}_{WG}\ , J = 9.$
(0.25, 0.1)	4.64123E-05	4.7799E-03
(0.5, 0.1)	6.23950E-05	3.8654E-03
(0.75, 0.1)	4.64123E-05	2.7143E-03
(0.25, 0.5)	1.28684E-04	6.5798E-02
(0.5, 0.5)	1.72999E-04	7.3478E-02
(0.75, 0.5)	1.28684E-04	1.8764E-02
(0.25, 1.0)	1.13949E-03	3.2385E-01
(0.5, 1.0)	1.63140E-03	5.7492E-01
(0.75, 1.0)	1.13949E-03	4.2954E-01

5. Conclusion

The main objective of this article is to develop two accurate methods to solve nonhomogeneous heat and wave equations. Bernoulli operational matrix with collocation method and wavelet with Galerkin method have reduced the problem into the linear algebraic system. Some illustrative problems are given to ensure the high efficiency of the proposed algorithms.

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Conflicts of Interest

The authors declare no conflicts of interest regarding the publication of this paper.

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