Three-Dimensional Stress Concentration Factor in Finite Width Plates with a Circular Hole

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ABSTRACT

The three-dimensional stress concentration factor (SCF) at the edge of elliptical and circular holes in infinite plates under remote tension has been extensively investigated considering the variations of plate thickness, hole dimensions and material properties, such as the Poisson’s coefficient. This study employs three dimensional finite element modeling to numerically investigate the effect of plate width on the behavior of the SCF across the thickness of linear elastic isotropic plates with a through-the-thickness circular hole under remote tension. The problem is governed by two geometric non-dimensional parameters, i.e., the plate half-width to hole radius \((W/r)\) and the plate thickness to hole radius \((B/r)\) ratios. It is shown that for thin plates the value of the SCF is nearly constant throughout the thickness for any plate width. As the plate thickness increases, the point of maximum SCF shifts from the plate middle plane and approaches the free surface. When the ratio of plate half-width to hole radius \((W/r)\) is greater than four, the maximum SCF was observed to approximate the theoretical value determined for infinite plates. When the plate width is reduced, the maximum SCF values significantly increase. A polynomial curve fitting was employed on the numerical results to generate empirical formulas for the maximum and surface SCFs as a function of \(W/r\) and \(B/r\). These equations can be applied, with reasonable accuracy, to practical problems of structural strength and fatigue, for instance.

Keywords: Stress Concentration Factor; Plates with Circular Hole; Finite Element Analysis

1. Introduction

Many applications in engineering employ components with a circular hole. In the specific case of perforated plates under cyclical load, the effect of stress concentration can propagate cracks and compromise their structural integrity. The stress concentration near a geometric discontinuity in a plate is frequently described by the stress concentration factor (K), defined as the ratio of the actual stress \((\sigma)\) acting on that region to the stress \((\sigma_0)\) applied to the plate extremity.

Howland [1] studied the stress around the central hole on finite width plates, using bipolar coordinates and biharmonic functions. The solution is iterative, so accuracy may be successively improved and the results were compared with photo-elastic experiments. Timoshenko and Goodier [2] and Muskhelishvili [3] presented classical solutions for bi-dimensional analysis of stress concentration along the hole edges on infinite width plates.

Based on the Theory of Linear Elasticity for plane strain or stress problems, Koiter [4] demonstrated that when the diameter of a centered, through-the-thickness hole approaches the plate width, the ratio between the maximum stress at the hole edge and the average stress in the reduced section \((\sigma_{max}/\sigma_{av})\) is in the limit equal to 2. Parks and Mendoza [5] employed experimental analysis with strain gages to study the behavior of plates when the ratio between the hole diameter and the plate width is equal to 0.98 and 0.99, and showed that \(\sigma_{max}/\sigma_{av}\) approaches 2. Wahl [6] used a simple non-linear formula, applied to small relations between the hole wall minimum thickness and the plate width and also showed that \(\sigma_{max}/\sigma_{av}\) tends to 2 and it approaches 1 as the load is increased. Cook [7] employed a geometrically non-linear finite element model to analyze elastic plates with Young’s Modulus and Poisson’s coefficient respectively equal to 100 GPa and zero, and varied the remotely applied stress from zero to 1000 MPa and the ratio between the hole...
wall minimum thickness and the plate width from $10^{-2}$ to $10^{-6}$. The results indicated that $\sigma_{\text{max}}/\sigma_{\text{av}}$ decreases from 1.94 to 1 as the load intensity is increased. Pradhan [8] employed a plane stress state finite element model for isotropic and composite plates and showed that the maximum stress at the edge of a centered, through-the-thickness hole strongly depends on the material property. The compilation of the stress concentration factors published by Pilkey [9] is usually considered an important reference for design.

Fatigue analysis methods for cracks and other forms of stress concentration are generally developed using bi-dimensional (2D) models. However, when these models are applied to certain types of problems where the geometry or material properties may lead to accentuated tri-dimensional (3D) stress concentrations the results can be inaccurate, as demonstrated by Bellett et al. [10] and Bellett and Taylor [11]. A variation of the SCF through the wall thickness of infinite width plate with elliptic and circular holes in isotropic material under remote tensile stress has been systematically investigated using finite element methods by Altenhof and Zamani [12], She and Guo [13,14], Yu et al. [15] and Yang et al. [16]. These analyses show that the maximum value of stress concentration occurs near but not on the plate surface as the thickness increases and this effect is more significant for elliptic holes. Kubair and Bhanu-Chandar [17] employed a FEM to investigate the effect of inhomogeneous material properties on the SCF for plates with a central circular hole subject to a remote stress. The material is functionally graded, that is, its properties vary spatially. A parametric study indicated that the SCF reduces when the Modulus of Young is progressively increased towards the center of the hole and that the angular position of maximum stress on the plate surface is unaffected by the material inhomogeneity. More recently, Chao et al. [18] presented an analytical solution for the stress field at an infinity plate with reinforced elliptical hole subjected to an inclined uniform remote tensile stress. The material properties for the reinforcement layer may differ from the plate properties, and the reinforcement layer is limited by two confocal ellipses. Rezaeepazhand and Jafari [19] carried out an analytical investigation for the stress distribution in isotropic plates with centered holes with different shapes, subjected to remote uniaxial tensile stress. In the parametric study circular, triangular, square and pentagon holes were considered, furthermore, the cut-out shape, bluntness and orientation were also taken into account. The analytical results were compared with FEM simulations and showed that the SCF for square holes are smaller than for a similar plate with circular holes and that triangular and pentagon holes yield higher SCFs.

The objective of this study is to evaluate the variation of the stress concentration factor through the thickness for isotropic plates, with through-the-thickness circular holes, subject to remote tensile stress and to investigate the effect of plate width on the results. A finite element model was elaborated with various widths and thicknesses to allow a comprehensive parametric evaluation of this variation.

2. Definition of Geometric Parameters

Figure 1 shows a plate with a through-the-thickness circular hole, submitted to uniaxial stress $(\sigma_x)$, where the geometric parameters used in this study can also be seen. The plate width, thickness and length are respectively equal to $2W$, $B$ and $2H$. The circular hole, located at the center of the plate, has a radius equal to $r$. The fixed geometric parameters are $r = 5$ mm and $H = 100$ mm, therefore the ratio $H/r$ is equal to 20, which ensures that the stress is applied far enough from the hole. The dimensionless variable ratios employed in the analyses are: $W/r = 1.2, 1.4, 1.6, 1.8, 2.0, 2.2, 2.5, 3, 4, 5, 6, 10$ and $20$ and $B/r = 0.2, 1, 2, 3, 4, 5, 6, 10, 20$ and $30$, totaling 130 simulation cases. The results in the analyses are related to these dimensionless parameters in such a way that an ample range of plate widths and thicknesses is investigated. Nevertheless, whenever convenient the results will be only presented for a smaller set of data.

3. Numerical Analysis

To study the stress concentration distribution along the edge of a circular hole and plate thickness, the finite element program ANSYS® [20] was used to obtain the solutions of the models with the geometric relations previously presented. The element SOLID45, defined by eight nodes with three degrees of freedom at each node (translations in the $x$, $y$ and $z$ nodal directions) was employed. This element is appropriate for the analysis of solid structures (3D) and recommended for linear applications. To reduce processing time, a model representing only $1/8$ of the plate was elaborated, taking advantage of

Figure 1. Dimensions of the plate with a circular hole and load.
the conditions of symmetry in the mid-planes, with the edge of the plate remaining free, as seen in Figure 2. In half-width \( W \), on mid-plane \( yz \), the node degrees of freedom in the \( x \) direction were restricted. In half-length \( H \), on mid-plane \( zx \), the node degrees of freedom in the \( y \) direction were also restricted. And, lastly, in half-thickness \( B/2 \), on mid-plane \( xy \), the node degrees of freedom in the \( z \) direction were restricted. The plate half-thickness was then divided in layers of elements, in such a way as to achieve a greater discretization in the region near the hole where a strong stress concentration needs to be captured, and elsewhere a regular mesh was devised ensuring that the aspect ratio of 1:20 in any element of the model was not violated. The elements along the half-thickness \( B/2 \) had a decreasing mesh, from the center to the free surface. Thus, the ratio of the nodal spacing \( \frac{1}{nBB} \) varied between 1 and 8, see detail in Figure 2. The discretization adopted was obtained through mesh convergence tests for models with \( Br = 2 \) and 10, with the number of degrees of freedom varying between 13,680 and 292,158. Hence, it was possible to ensure the same mesh spacing around the hole edge for any values of \( Wr \) and \( Br \), with an error inferior to 1% in the values of maximum stress.

A longitudinal load was applied in the \( y \) direction through a tension load on the surface of the plate extremity, ensuring that stress did not exceed the yield stress for ordinary naval steel based on the ASTM A131 [21] standard.

The model was then considered geometrically and physically linear, bearing in mind that small displacements and an elastic linear regime were assumed. To represent the material parameters applied in the model, the Young modulus and the Poisson’s coefficient for steel were defined as 200 GPa and 0.333, respectively.

4. Results

Timoshenko and Goodier [2] employed the theory of elasticity for problems in two-dimension, using polar coordinates to demonstrate that the distribution of radial, shear and circumferential stresses in the area of a circular hole in the center of an infinite width plate, as shown in Figure 3, is given via the following expressions:

\[
\sigma_r = \tau_{r\theta} = 0
\]

\[
\sigma_{\theta} = \sigma_y (1 - 2 \cos 2\theta)
\]

It is clear that the maximum value of \( \sigma_{\theta} \) occurs on the plane of the transversal section to the plate passing through the center, i.e., when \( \theta = \pi/2 \) or \( 3\pi/2 \). Thus, the maximum stress is three times greater than the uniform tension \( \sigma_y \) applied in the plate extremities, in other words, the theoretical stress concentration factor \( K_o \) is:

\[
K_o = \frac{\sigma_{\theta,max}}{\sigma_y} = 3
\]

In Figures 4(a)-(c) the curves for various half-width to hole radius relations \( (W/r) \) are presented, which show the variation of the stress concentration factor at the hole edge \( (K/K_o) \) along the plate thickness. The distance from the plate middle plane to the considered point is represented by the dimensionless parameter \( 2d/B \), where \( d \) is this distance and \( B/2 \) the half-thickness, hence \( 0 \leq 2d/B \leq 1 \). Figures 4(a)-(c) show the curves for the ratios of thickness to the hole radius \( B/r = 0.2, 4 \) and 30, respectively. The curves for \( B/r = 1, 2, 3, 5, 6 \) and 10 and \( B/r = 20 \) are respectively similar to \( B/r = 0.2, 4 \) and 30, hence to save space they are not shown here. Due to the overlay of the curves and the consequent line congestion on the bottom part of Figure 4, the curves corresponding to \( W/r = 1.8, 2.2, 2.5, 3, 5, 6, 10 \) are likewise not shown.

The general interpretation for the behavior of the curves shown in Figures 4(a)-(c) indicates a significant increase in \( K/K_o \) with a decrease in \( W/r \) for all cases, especially for the range of \( W/r = 1.2 \) to 4. Fur-

Figure 2. Constructive detail of the numerical model.

Figure 3. Plate with a centered hole under remote tension.
thermore, the maximum stress concentration factor for $W/r = 1.2$ is always greater than four times the theoretical 2D value for any of the $B/r$ relations. In Figure 4(a), when $B/r = 0.2$, the curves present profiles of stress concentration factors almost constant with thickness for the different values of $W/r$. The curves of Figures 4(b) and (c), corresponding to $B/r = 4$ and 30, respectively, demonstrate an almost constant function up to a specific distance from the middle plane, and from this point there is an accentuated decrease of $K/K_o$ towards the plate surface.

Figure 5 indicates a variation of the ratio of the distance to the middle plane ($d_{max}$) from the point of maximum stress concentration factor to the half-thickness of the plate ($B/2$), as a function of the ratio of the plate half-width to the hole radius ($W/r$) with varying $B/r$ ratios.

For $B/r \leq 3$, the maximum stress concentration factors ($K_{max}$) remain at the plate middle plane, that is, $2d_{max}/B = 0$, for any $W/r$ value. When the ratio of thickness to the hole radius is greater than 3 ($B/r > 3$) and the ratio plate half-width to the hole radius is greater than 2 ($W/r > 2$), the location of maximum stress concentration factor displaces from the middle plane and approaches, but it never reaches, the plate free surface. Besides, for every $B/r$ curve, beginning at a specific $W/r$ value, the position of the point of $K_{max}$ remains constant.

Figures 6(a) and (b) respectively demonstrate the ratios of maximum stress concentration ($K_{max}$) and the surface stress concentration ($K_s$) to the 2D theoretical value ($K_o$), as a function of the $B/r$ ratio, when the ratio of plate half-width to the hole radius ($W/r$) is varied. To avoid the accumulation of curves in the lower part of the graph and, thus, lose clear sight, the curves referring to $W/r = 4, 5, 6$ and 10 are not shown.
and surface stress concentration factors \((K_{\text{max}}/K_r)\) as a function of the \(B/r\) ratio, for various \(W/r\) ratios. To avoid overlaying, the curves for \(4 \leq W/r \leq 10\) relations are not displayed. Figure 7 indicates that the values for the \(K_{\text{max}}/K_r\) ratios are very close to one only for \(B/r = 0.2\); and when the plate thickness is increased the ratio \(K_{\text{max}}/K_r\) values strongly decrease with \(B/r\). When \(B/r > 10\), except for \(W/r = 1.2\), the \(K_{\text{max}}/K_r\) curves asymptotically converge to values close to 1.24.

Figures 8(a) and (b) indicate, respectively, \(K_{\text{max}}/K_o\) and \(K_r/K_o\) ratios, as a function of \(W/r\), for various \(B/r\) ratios. The curve obtained by Howland [1] is also included in this graph. The curves \(K_{\text{max}}/K_o\) and \(K_r/K_o\) in the range \(4 \leq B/r \leq 30\) practically blend together, so to improve visualization it was decided to trace only the curves for the ends of this interval.

When comparing the numeric values obtained for thin plates, typically represented herein by \(B/r = 0.2\), with Howland’s [1] results, some small discrepancies are observed between the curves for \(K_{\text{max}}/K_o\), and better adherence for \(K_r/K_o\). This is expected because photoelastic measurements, in which Howland’s [1] equations were confronted, had been performed on plate surfaces.

It becomes evident in Figures 8(a) and (b) that Howland’s [1] results are valid only for plates with the \(B/r\) relation close to 0.2. Figure 8(a) indicates that Howland’s [1] curve underestimates \(K_{\text{max}}/K_o\) for other \(B/r\) relations. Figure 8(b) demonstrates that Howland’s [1] curve overestimates the \(K_r/K_o\) values when \(B/r > 0.2\).

5. Design Formulas

Due to the non-linear characteristics of the variables in-

![Figure 6. Variation of \(K_{\text{max}}/K_o\) and \(K_r/K_o\) as a function of \(B/r\). (a) Values for \(K_{\text{max}}/K_o\) for several values of \(W/r\); (b) Values for \(K_r/K_o\) for several values of \(W/r\).](image)

![Figure 7. Variation of \(K_{\text{max}}/K_r\) as a function of \(B/r\) for several \(W/r\) relations.](image)
Figure 8. Variation of $K_{\text{max}}/K_0$ and $K_s/K_0$ as a function of $W/r$. (a) Values of $K_{\text{max}}/K_0$ for several $B/r$ values; (b) Values of $K_s/K_0$ for several $B/r$ values.

The empirical formula for the $K_{\text{max}}/K_0$ curves shown in Figure 8(a) can be expressed by:

$$K_{\text{max}}/K_0 = a_3 \beta^3 + a_2 \beta^2 + a_1 \beta + a_0$$

where:

- $a_0 = -0.0064 \alpha^2 + 0.0122 \alpha + 1.2492$
- $a_1 = 0.0026 \alpha^2 + 0.0001 \alpha - 0.4951$
- $a_2 = -0.0022 \alpha^2 + 0.0031 \alpha + 0.4562$
- $a_3 = 0.0007 \alpha^2 - 0.0024 \alpha - 0.1580$

With parameters $\alpha$ and $\beta$ respectively given by

$$\alpha = \ln (B/r) \quad \text{and} \quad \beta = \ln \left[ \ln (W/r) \right].$$

For $K_s/K_0$ displayed in Figure 8(b) the empirical formula is expressed by:

$$K_s/K_0 = b_3 \beta^3 + b_2 \beta^2 + b_1 \beta + b_0$$

where:

- $b_0 = 0.004 \alpha^3 - 0.004 \alpha^2 - 0.0691 \alpha + 1.1191$
- $b_1 = 0.0012 \alpha^3 - 0.0032 \alpha^2 + 0.0148 \alpha - 0.4424$
- $b_2 = 0.0018 \alpha^3 - 0.0057 \alpha^2 - 0.0213 \alpha + 0.4307$
- $b_3 = -0.0022 \alpha^3 + 0.0058 \alpha^2 + 0.0137 \alpha - 0.1475$

The maximum errors of the values obtained by Equations (4) and (5), when compared with the $K_{\text{max}}/K_0$ and $K_s/K_0$ corresponding values obtained by finite element analyses are less than 1.50% and 1.68%, respectively.

6. Conclusions

The present study presents the variation of the stress concentration factor through the thickness of linear elastic isotropic plate with through-the-thickness circular hole, subject to remote tensile stress and investigates the effect of plate width on the behaviour of the results. A finite element method was used to analyze the 3D plate solid structure, using an eight nodes element with three degrees of freedom per node, which is appropriate and recommended for use in linear applications. A strong influence of the ratio plate half-width to hole radius ($W/r$) on the stress concentration factor was observed during the analyses. The study demonstrated that with a fixed thickness to hole radius $B/r$ relation the behavior of the stress concentration factors presented similar characteristics for different $W/r$ relations, but at rather different $K/K_0$ levels.

The largest $K_{\text{max}}$ observed for the set of geometric parameters investigated in this paper corresponds to the plate with $1.2 \leq W/r \leq 20$ and plate thickness to hole radius of $0.2 \leq B/r \leq 30$.

For $K_{\text{max}}/K_0$ curves versus $W/r$, (a) Values of $K_{\text{max}}/K_0$ for several $B/r$ values; (b) Values of $K_s/K_0$ for several $B/r$ values.

The largest $K_{\text{max}}$ observed for the set of geometric parameters investigated in this paper corresponds to the plate with $W/r=1.2$ and $B/r=2.0$, and is equal to 4.23 times the stress concentration factor for the infinite width plate ($K_0$). The largest surface stress concentration factor ($K_s/K_0$) also occurs in the plate with $W/r=1.2$ and $B/r=2.0$, and is equal to 3.99. For values of the ratio $B/r \leq 3.0$, the point where $K_{\text{max}}$ occurs is independent of $W/r$ and is always on the plate middle surface. For plates with $B/r > 3.0$ the distance from the point of $K_{\text{max}}$ to the mid-plane ($d_{\text{max}}$) depends on the $W/r$ and $B/r$ ratios. When $B/r > 10$ the value of the maximum stress concentration factor ($K_{\text{max}}$) approximately tends to 24% greater than the surface stress concentration factor ($K_s/K_0$) for $W/r > 1.2$. It was demonstrated that the $K_{\text{max}}/K_0$ curves versus $W/r$, in the same way that the $K_s/K_0$ curves versus $B/r$, for all $W/r$ ratios, present asymptotic behavior for $W/r < 0.2$. The $K_{\text{max}}/K_0$ curve versus $W/r$ numerically developed when $B/r = 0.2$ presents small...
discrepancies when compared to Howland’s. For all other \( B/r \) relations the \( K_{max}/K_0 \) are underestimated by Howland’s values. The \( K_s/K_0 \) curve versus \( W/r \) developed numerically when \( B/r = 0.2 \) presents good consonance with Howland’s prediction. However, for all other \( B/r \) relations the \( K_s/K_0 \) are overestimated by Howland. The empirical formulas developed from the numerical results presented maximum errors for \( K_{max}/K_0 \) and \( K_s/K_0 \) respectively equal to 1.50% and 1.68%.

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