Simulation and Modelling of Smart Beams with Robust Control Subjected to Wind Induced Vibration

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Received May 22, 2012; revised June 28, 2012; accepted July 9, 2012

ABSTRACT
This paper presents some recent developments in modelling and numerical analysis of piezoelectric systems and controlled smart structures based on a finite element formulation with embedded control. The control aims at vibration suppression of the structure subjected to external disturbances, like wind and noise, under the presence of model inaccuracies, using the available measurements and controls. A smart structure under dynamic loads is analysed and comparison between results for beam with and without control is made. The numerical results show that the control strategy is very effective and suppresses the vibrations of the structure.

Keywords: Smart Structures; Finite Element Formulation; Uncertainty; Robust Performance; Reduced Order Control

1. Introduction
The intelligent material sector has broadly motivated research interest during the last decades. Materials are classified as intelligent if they exhibit the ability to sense a certain stimulus and initiate a reaction to it in a controlled manner [1-11]. An intelligent structure is a structural system or member with embedded sensors and actuators, combined with a control system that automates the structure’s reaction to external loadings acting on it, by attenuating their influence or improving desirable responses. Development in this scientific discipline is aided by advances in material science as well as in the field of control. Material science research has spurred the development of new materials that can be incorporated in structures to be used as actuators and sensors [9-11].

In this paper, we address the problem of vibrations of intelligent structures. Stimuli may come from external perturbations of the system, disturbances or excitation that may cause structural vibrations, such as wind loading or earthquakes. An intelligent structure is expected to be able to sense the vibration and counteract it in a controlled fashion, so that vibration of the system be reduced and contained. To that end, a number of intelligent materials may be used as actuators and sensors. [9-11] Piezoelectric materials, memory materials, electrostrictive and magnetostrictive materials are such materials. In this work, we focus on the use of piezoelectric materials, given that they exhibit good sensing and actuation properties.

Among the commonly used optimal control schemes are LQR, and \( H_\infty \). It is known that if the controller is not robust enough, the uncertainties of the system may destroy the efficiency of the controller. The \( H_\infty \) control provides better robustness than LQR [12-14]. The aim of this work is to design an \( H_\infty \) robust controller for a beam bonded with piezoelectric sensors and actuators and to investigate the behaviour of the controlled beam. First, a detailed shear-deformable (Timoshenko) model for a laminated beam structure is developed. A finite element formulation is presented for the model. Quadratic Hermitian polynomials are used for the transverse and rotational displacements, respectively. The differential equations are based on the Timoshenko beam theory [15]. The governing state equation is established and used for the design of the control. The numerical simulations carried out on the laminated beam shows that the vibration of the system is significantly suppressed within the permitted actuator voltages. Herein the integration of control into a home-made finite element code developed in MATLAB is presented [16]. The numerical solution of the \( H_\infty \) feedback controller has been done by using a nonconvex, non differentiable optimization approach with the usage of HIFOO software within MATLAB [14,17]. For the numerical results wind type loads are taken into account. The effectiveness of the technique in the modelling of the standard uncertainties is also presented.

2. Mathematical Modelling
A cantilever slender beam with rectangular cross-sec-
tions is considered. Four pairs of piezoelectric patches are embedded symmetrically at the top and the bottom surfaces of the beam, as shown in Figure 1. The beam is made from graphite-epoxy T300-976 and the piezoelectric patches are PZT G1195N. The top patches act like sensors and the bottom like actuators. The resulting composite beam is modelled by means of the classical laminated technical theory of bending. Let us assume that the mechanical properties of both the piezoelectric material and the host beam are independent in time. The thermal effects are considered to be negligible as well [5, 6].

The beam has length \( L \), width \( W \) and thickness \( h \). The sensors and the actuators have width \( b_s \) and \( b_a \) and thickness \( h_s \) and \( h_a \), respectively. The electromechanical parameters of the beam are given in the Table 1.

### 2.1. Reduced Piezoelectric Equations

In order to derive the basic equations for piezoelectric sensors and actuators (S/As), we assume that:
- The piezoelectric S/A are bonded perfectly on the host beam;
- The piezoelectric layers are much thinner than the host beam;
- The piezoelectric material is homogeneous, transversely isotropic and linearly elastic;
- The piezoelectric S/A are transversely polarized (in the z-direction) [5].

Under these assumptions the three-dimensional linear constitutive equations are given by [4],

\[
\begin{bmatrix}
\sigma_{xx} \\
\sigma_{zz}
\end{bmatrix} = \begin{bmatrix}
Q_{11} & 0 \\
0 & Q_{55}
\end{bmatrix} \begin{bmatrix}
\varepsilon_{xx} \\
\varepsilon_{zz}
\end{bmatrix} - \frac{d_{31}}{\varepsilon_3} E_z
\]

(1)

\[
D_z = Q_{11} d_{31} \varepsilon_{xx} + \varepsilon_3 E_z
\]

(2)

where \( \sigma_{xx} \), \( \sigma_{zz} \) denote the axial and shear stress components, \( D_z \) denotes the transverse electrical displacement; \( \varepsilon_{xx} \) and \( \varepsilon_{zz} \) are a axial and shear strain components; \( Q_{11}, Q_{55} \) denote elastic constants; \( d_{31} \), and \( \varepsilon_3 \), denote piezoelectric and dielectric constants, respectively. Equation (1) describes the inverse piezoelectric effect and Equation (2) describes the direct piezoelectric effect. \( E_z \) is the transverse component of the electric field that is assumed to be constant for the piezoelectric layers and its components in xy-plain are supposed to vanish. If no electric field is applied in the sensor layer, the direct piezoelectric Equation (2) gets the form,

\[
D_z = Q_{11} d_{31} \varepsilon_{xx}
\]

(3)

and it is used to calculate the output charge created by the strains in the beam [3].

### 2.2. Equations of Motion

The length, width and thickness of the host beam are denoted by \( L, b \) and \( h \), respectively. The thickness of the sensor and actuator is denoted by \( h_s \) and \( h_a \). We assume that:
- The beam centroidal and elastic axis coincides with the x-coordinate axis so that no bending-torsion coupling appears;
- The axial vibration of the host beam centreline is considered negligible;
- The displacement field \( \{u\} = (u_x, u_y, u_z) \) is obtained based on the usual Timoshenko assumptions,

\[
\begin{align*}
\omega(x, y, z) &\approx \varphi(x, t) \\
u_y(x, y, z) &\approx 0 \\
u_z(x, y, z) &\approx \omega(x, t)
\end{align*}
\]

(4)

where \( \varphi \) is the rotation of the beams cross-section about the positive y-axis and \( \omega \) is the transverse displacement of a point of the centroidal axis \( (y = z = 0) \).

The strain displacement relations can be applied to Equation (4) to give,

\[
\varepsilon_{xx} = z \frac{\partial \varphi}{\partial x} \\
\varepsilon_{yy} = z \frac{\partial \varphi}{\partial x}
\]

(5)
We suppose that the transverse shear deformation is equal to zero [2]. In order to derive the equations of the motion of the beam we use Hamilton’s principle,

$$\int_{t_1}^{t_2} (\delta T - \delta U + \delta W) \, dt = 0$$

(6)

where $T$ is the total kinetic energy of the system, $U$ is the potential (strain) energy and $W$ is the virtual work done by the external mechanical and electrical loads and moments. The first variation of the kinetic energy is given by,

$$\delta T = \frac{1}{2} \int_{V} \rho \left( \frac{\partial u}{\partial t} \right)^T \left( \frac{\partial u}{\partial t} \right) \, dV$$

(7)

$$= \frac{b}{2} \int_{0}^{h} \int_{-h/2}^{h/2} \rho \left( \begin{array}{c} \frac{\partial u}{\partial x} \\ \frac{\partial u}{\partial z} \\ \frac{\partial \omega}{\partial x} \\ \frac{\partial \omega}{\partial z} \end{array} \right) \left( \begin{array}{c} \frac{\partial u}{\partial x} \\ \frac{\partial u}{\partial z} \\ \frac{\partial \omega}{\partial x} \\ \frac{\partial \omega}{\partial z} \end{array} \right) \, dz \, dx$$

The first variation of the kinetic energy is given by,

$$\delta U = \frac{1}{2} \int_{V} \delta \left[ \frac{\partial \epsilon}{\partial x} \right]^T \left[ \sigma \right] \, dV$$

(8)

$$= \frac{b}{2} \int_{0}^{h} \int_{-h/2}^{h/2} \rho \left( \begin{array}{c} \frac{\partial \epsilon}{\partial x} \\ \frac{\partial \epsilon}{\partial z} \\ \frac{\partial \gamma}{\partial x} \\ \frac{\partial \gamma}{\partial z} \end{array} \right) \left( \begin{array}{c} \frac{\partial \epsilon}{\partial x} \\ \frac{\partial \epsilon}{\partial z} \\ \frac{\partial \gamma}{\partial x} \\ \frac{\partial \gamma}{\partial z} \end{array} \right) \, dz \, dx$$

If the load consists only of moments induced by piezoelectric actuators and since the structure has no bending twisting couple then the first variation of the work has the form [6],

$$\delta W = \frac{b}{2} \int_{0}^{h} M^4 \left( \frac{\partial \phi}{\partial x} \right) \, dx$$

(9)

where $M^4$ is the moment per unit length induced by the actuator layer and is given by,

$$M^4 = \frac{h}{2} \int_{-h/2}^{h/2} \frac{\partial \phi}{\partial x} \, dz = \frac{b}{2} \int_{-h/2}^{h/2} E_z \, dz$$

(10)

2.3. Finite Element Formulation

We consider a beam element of length $L_e$, which has two mechanical degrees of freedom at each node: one translational $\omega_1$ (respectively $\omega_2$) in direction $z$ and one rotational $\phi_1$ (respectively $\phi_2$), as it is shown in Figure 2. The vector of nodal displacements and rotations $q_e$ is defined as [3],

$$q_e = [\omega_1, \phi_1, \omega_2, \phi_2]$$

(11)

The beam element transverse deflection $\omega(x, t)$ and the beam element rotation $\phi(x, t)$ of the beam are continuous and they are interpolated within by Hermitian linear shape functions $H^H$ and $H^R$ as follows [4],

$$\omega(x, t) = \sum_{i=1}^{4} H^H(x) q_i(t)$$

$$\phi(x, t) = \sum_{i=1}^{4} H^R(x) q_i(t)$$

(12)

This classical finite element procedure leads to the approximate (discretized) variation problem. For a finite element the discrete differential equations are obtained by substituting the discretized expressions 12 into Equations (7) and (8) to evaluate the kinetic and strain energies. Integrating over spatial domains and using the Hamiltons principle 6 the equation of motion for a beam element are expressed in terms of nodal variable $q$ as follows,

$$M \ddot{q}(t) + D \dot{q}(t) + K q(t) = f_m(t) + f_e(t)$$

(13)

where $M$ is the generalized mass matrix, $D$ the viscous damping matrix, $K$ the generalized stiffness matrix, $f_m$ the external loading vector and $f_e$ the generalized control force vector produced by electromechanical coupling effects. The independent variable $q(t)$ is composed of transversal deflections $\omega$ and rotations $\phi$, i.e., [7]

$$q(t) = [\omega_1, \phi_1, \omega_2, \phi_2]$$

(14)

where $n$ is the number of nodes used in analysis. Vectors $\omega$ and $f_m$ are positive upwards. To transform to state-space control representation, let (in the usual manner),

$$\dot{x}(t) = \begin{bmatrix} \dot{\omega}_1 \\ \dot{\phi}_1 \\ \dot{\omega}_2 \\ \dot{\phi}_2 \end{bmatrix}$$

(15)

Furthermore to express $f_e(t)$ as $Bu(t)$ we write it as $f_e u$ where $f_e$ the piezoelectric force is for a unit applied on the corresponding actuator, and $u$ represents the voltages on the actuators. Furthermore, $d(t) = f_m(t)$ is the disturbance vector [18]. Then,

$$\dot{x}(t) = \begin{bmatrix} O_{2n \times 2n} & I_{2n \times 2n} \\ -M^{-1}K & -M^{-1}D \end{bmatrix} x(t) + \begin{bmatrix} O_{2n \times 2n} \\ M^{-1}f_e \end{bmatrix} u(t) + \begin{bmatrix} O_{2n \times 2n} \\ M^{-1} \end{bmatrix} d(t)$$

(16)

$$= Ax(t) + Bu(t) + Gd(t)$$

$$= Ax(t) + B(\bar{u}(t))$$

(17)
The previous description of the dynamical system will be augmented with the output equation (some displacements or velocities are measured) [4],
\[ y(t) = \begin{bmatrix} x_1(t) & x_2(t) & \cdots & x_{n-1}(t) \end{bmatrix}^T = Cx(t) \] (18)

In this formulation \( u \) is \( n \times 1 \) (at most, but can be smaller), while \( d \) is \( 2n \times 1 \). The units used are compatible for instance \( m \), rad, sec and \( N \).

3. Statement of the Robust Control Problem

The optimal control problem is initially studied for the nominal system, i.e., the beam with known elastic, piezoelectric and viscous properties. A more realistic question concerning the robustness of the control in the presence of defects is also addressed. The fact that the system is in fact subject to disturbances, such as the wind power, as well as the external additional influences, is taken into account. The mathematical pattern being used in the design is an approximation of the real one. Further, two control laws for the composite beam are designed in order to suppress the vibrations. Because of its linearity and easy implementation, the linear quadratic regulator (LQR) [13] is presented first. The response of the controlled nominal and approximated system will be discussed later on.

A system analysis is presented in order to check whether the control in the presence of defects is also addressed. The fact that the system is not accurately known and the input equation (some disturbances are measured) [4], is the structure singular value of \( M \) given the uncertainty set \( B_{\Delta} \).

2) The system \((M, \Delta)\) exhibits robust performance if,
\[ \sup_{\Delta \in B_{\Delta}} \mu(M(j\omega)) < 1 \] (21)

where,
\[ \Delta_u = \begin{bmatrix} \Delta_p & 0 \\ 0 & \Delta \end{bmatrix} \]

and \( \Delta_p \) has the same structure as \( \Delta \) but dimensions corresponding to \((w,z)\). Unfortunately, only bounds on \( \mu \) can be estimated.

To proceed let us assume uncertainty in the \( M, D \) and \( K \) matrices of the form,
\[ M = M_0 (I + m_p \delta_m) \]
\[ D = D_0 (I + d_p \delta_D) \]
\[ K = K_0 (I + k_p \delta_k) \] (23)

with,
\[ \|\Delta\|_\infty = \begin{bmatrix} \delta_M \\ \delta_D \\ \delta_K \end{bmatrix} \leq 1 \] (24)

This means that we are allowing a percentage deviation from the nominal values [4]. With these definitions Equation (13) becomes,
\[ M_0 \ddot{q}(t) + D_0 \dot{q}(t) + K_0 q(t) \]
\[ = \dot{\bar{D}}_D q(t) + m \dot{f}_m(t) + f_c(t) \] (25)

where,

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aim is to study the response of the composite beam in the presence of defects and damages.

Two kinds of dynamic loading are used as disturbances:

- Periodic sinusoidal loading pressure acting on the side of the structure simulating a strong wind. A sinusoidal load with an amplitude of 10 N and frequency of 6.2832 rad/sec, has been considered.
- A typical stochastic wind-type load on the side of the structure (Figure 4).

4. Numerical Examples

For the numerical simulations a cantilever composite beam with viscous and piezoelectric layers bonded on its top and bottom and discretized with four finite elements, is used. A finer finite element discretization, which certainly is required for the approximation of higher frequencies, does not change the trend of the results. Our

4.1. Results

The first load analyzed is sinusoidal load acting on the side of the structure. Figure 5 shows the dynamical response for the displacements of the uncontrolled and

\[
\dot{x}(t) = Ax(t) + Bu(t) + Gd(t) + G_q q_u(t)
\]

In this way we treat uncertainty in the original matrices as an extra uncertainty term.

![Figure 4. Corresponding wind load acting.](image)

![Figure 5. Response of the four nodes for the displacement with LQR and $H_\infty$ control and response without control.](image)
A controlled beam with LQR control [12,13] and $H_\infty$ control [14,19], for the four nodes of the beam. Figure 6 shows the dynamical response for the rotations of the uncontrolled and controlled beam with LQR control and $H_\infty$ control strategy, for the four nodes of the beam. The beam with $H_\infty$ control keeps in equilibrium and we have zero displacements, complete vibration reduction is achieved. The comparison of the open and closed loop frequency response of the system are shown in Figure 7, as shown in figure, there is a significant improvement in the effect of disturbance on error up to the frequency of 1000 Hz. Figure 8 shows the control voltages for the four nodes of the beam. The control voltages for the disturbance rejections of the beam are less than 500 volt.

Comparison with the open loop response for the same plant shows the good performance of the $H_\infty$ controller. Results are very good, and the beam remains in equilibrium. Reduction of vibrations is observed, while piezoelectric add-ons produce voltage within their tolerance limits (±500 volt).

Then a typical wind load (Figure 4) acting on the side of the structure. The wind load is a real life wind speed measurements in relevance with time that took place in Estavromenos of Heraklion Crete. We transform the wind speed in wind pressure with,

![Graphs showing rotation response for different nodes with and without control](image1)

**Figure 6.** Response of the four nodes on the vibrating beam for the rotation with LQR and $H_\infty$ control and without control.

![Graph showing singular value for $H_\infty$ control strategy](image2)

**Figure 7.** Singular value for $H_\infty$ control strategy.
where \( V = \text{velocity}, \rho = \text{density} \) and \( C_v = 1.5 \).

In Figure 9 the structural displacement responses are reported, for the four nodes of the beam. It is possible to see clearly the benefit induced by the control on the maximum value of these displacements. \( H_\infty \) controller results are very satisfactory and prove that \( H_\infty \) control can reduce smart structures vibrations. In Figure 10 we can see the structural rotations for the four nodes of the beam. The beam with \( H_\infty \) control keeps in equilibrium and we have almost zero displacements.

In Figure 11 we can see the control profile for the four nodes of the beams. As we can see the voltage is more less than 500 Volt, which is the piezoelectric limit.

### 4.2. Reduced Order Control

The \( H_\infty \) controller found is of order 24. The fact that controller order, which is equal to the order of the system, is relatively higher than the order of classical controllers such as LQR has led a number of researchers to develop order reduction algorithms. The most widely used such algorithm, known as HIFOO, has been implemented in a Matlab environment, and is the one used in the following procedure [17, 20].

The general problem is to compute a controller of reduced order \( n < 24 \) while retaining the performance of the \( H_\infty \) criterion as well as the behavior of a full order controller for the given system [17, 20].

Nonsmooth variation analysis and related computational methods are powerful tools that can be effectively applied to identify local minimizers of nonconvex optimization problems arising in fixed-order controller design. Our computational methods found a 24th order controller that stabilized the system. Using the Matlab package HIFOO we can reduce this controller and stabilize the system with a 2nd order controller without difficu-
Figure 10. Rotations of the four nodes of the beam without and with $H_\infty$ control.

Figure 11. Control voltages for the four nodes of the beam using $H_\infty$ control.

culty, suggesting explicit formulas for the controller and for the closed loop system. Furthermore, our analytical techniques prove that this controller is locally optimal in the sense that there is no nearby controller with the same order for which the closed loop system has all its poles further left in the complex plane [20].

These approaches can be extended in order to take into account other key quantities of great practical interest, such as optimization of $H_\infty$ performance. In particular, with the help of a MATLAB toolbox called HIFOO ($H_\infty$ Fixed Order Optimization) we can reduce the order of the controller and have very good results. Figure 12, shows the response of the uncontrolled and control beam of the free end of the beam, using hifoo controller (a second order controller). Figure 13 shows the control voltages for the four nodes of the beam [20]. As we can see we need less energy then the 24th controller for reduced vibrations.

5. Conclusion

A mathematical formulation and finite element model for the vibration suppression of a cantilever beam with piezoelectric laminated surface and viscous layers and elastic core is presented in this paper. The design of the piezoelectric active control using LQR strategy and $H_\infty$ control theory for the nominal and damaged sandwich beam has been studied. The numerical results show that the $H_\infty$ control strategy is very effective and suppresses the vibrations of the beam. The vector of active control forces
subjected to $H_\infty$ performance criterion and satisfying the systems dynamic equations such that to reduce the wind excitations is determined. High robust performance and robust stability are achieved. This work clearly demonstrates the advantages of using advanced robust control theory for the design of practical smart structure.

6. Acknowledgements

The authors would like to thank the Wind Energy and Synthesis of Energy Systems Laboratory, of the Technological Educational Institute of Crete, for their offer on real life wind speed measurements in relevance with time that took place in Eustratomenos of Heraklion Crete.

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