Ferrimagnetism and Ferroelectricity of the Composite Matrix: SrBi$_2$Nb$_2$O$_9$(SBN)$_X$-BaFe$_{12}$O$_{19}$(BFO)$_{100-X}$

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ABSTRACT

A study of the dielectric, magnetic and structural properties of composites based on M-type barium hexaferrite BFO (BaFe$_{12}$O$_{19}$) and SBN (SrBi$_2$Nb$_2$O$_9$) is presented. The magneto-dielectric matrix composite (SrBi$_2$Nb$_2$O$_9$)$_x$(BaFe$_{12}$O$_{19}$)$_{100-x}$, ($x = 0, 25, 50, 75$ and $100$ wt$\%$) were prepared by using a new procedure based in the solid state method. X-ray powder diffraction patterns, Raman and Infrared spectroscopy, Mössbauer spectroscopy and scanning electron microscopy (SEM) were carried out for better understanding of the microstructural, dielectric and magnetic properties. Radiofrequency (RF) dielectric permittivity, dielectric loss measurements and magnetic and electric hysteresis loops properties are also discussed throughout this paper. The hysteresis loops showed that composite samples preserve the ferrimagnetism and ferroelectricity for hexaferrite when SBN is added to the composite (BFO$_{25}$P), although they become less coercive. In addition, the effects of organic binders group (TEOS, PVA and glycerin) on structural properties were also investigated.

Keywords: Composites; Magneto-Dielectric Materials

1. Introduction

The ceramics composite matrix are uniform, multiphases materials that have been strongly studied in recent works because of the possibility to obtain materials with desirable properties, especially for application in electronics devices. This properties wouldn’t be easily observed for a single phase material, requesting extreme conditions, for example, low temperatures, to its observation thus limiting its applications.

In this work, our main goal is to develop a dielectric material that is able to respond to both electric and magnetic stimulus, i.e. that is ferroelectric and ferromagnetic. To do so, we use the Aurivillius ceramic SrBi$_2$Nb$_2$O$_9$ and the Hexaferrite BaFe$_{12}$O$_{19}$. Such a material could be applied in the same way that common dielectrics (as dielectric resonator antennas, for example) but opening a wide range of possibilities to make the application of ceramics to electronic devices, memories and telecommunications more useful and powerful.

The Aurivillius family of compounds that can be represented by the general formula $[[\text{Bi}_2\text{O}_2]^{2+}(\text{A}_{m-1}\text{B}_m\text{O}_{3m+1})^{2-}]$, in which A can be a monovalent, divalent or trivalent cation or a combination of those in suitable proportions, B can be a tetravalent or pentavalent cation with m having values of 2, 3, 4... etc. Their structures comprise a stacking of n perovskite units of normal composition $\text{A}_{m-1}\text{B}_m\text{O}_{3m+1}$ interleaved with Bi$_2$O$_3$ layers along the pseudotetragonal c-axis known as bismuth layered ceramics (BLFC).

They are potential materials for microelectronic applications when desirable properties are reached such as high dielectric permittivity, high Curie temperature, high fatigue resistance [1-3]. However, dielectric loss is still too high for practical applications for electronic devices. Fortunately, the structure of perovskite allows easy replacement of cations and combinations among phases as well.

Recently, some researchers reported additions of impurities in bismuth-layered composites [1-3]. This work focused on a new alternative for a composite matrix by using barium hexaferrite.
M-type ferrite has a hexagonal structure and has been largely used for ferrofluids and magnetic devices. It has some interesting properties such as high saturation magnetization and high Curie temperature and has attracted interest for many years. By combining ferroelectric SBN ceramics and M-type ferrite, it was possible to produce the magneto-dielectric composite SrBi₂Nb₂O₉(SBN)ₓ-BaFe₁₂O₁₉(BFO)₁₀₀₋ₓ. By bringing these two materials together, it is expected to develop a bi-phase material, which would be ferroelectric and ferromagnetic. To do so, it would require SBN not to interact with BFO changing their characteristic properties, and the original phases can remain in the composite. The microstructure investigation is able to reveal whether the materials interact with each other forming new phases or not. If there is no undesirable chemical interactions between SBN and BFO, the new bi-phase ceramic will be formed and it will probably be ferroelectric and ferromagnetic.

The behavior of this new class of ceramics is not clearly known so far. Considering such context a precise investigation should be carried out for each composite combination. These matrix composites have interactions with magnetic and electric fields which are not completely understood because of cumbersome procedures for getting both ferroelectric and ferromagnetic properties under simple physical conditions. By using a composite matrix it is possible to mitigate those drawbacks and two non-interacting ferromagnetic and ferroelectric ceramics would provide magneto-dielectric properties. In this work, we have studied the microstructure, dielectric and magnetic properties of SrBi₂Nb₂O₉(SBN)ₓ-BaFe₁₂O₁₉ (BFO)₁₀₀₋ₓ matrix.

2. Experimental Procedure

SBN (SrBi₂Nb₂O₉) were synthesized from stoichiometric quantities of Bi₂O₃ (99.99%), Nb₂O₅ (99.99%), SrCO₃ (99.99%) and BFO (BaFe₁₂O₁₉) from BaO (99.99%) and Fe₂O₃ (99.99%) derived from solid-state reaction method. Both reactions of the uncalcined powders are shown below:

\[
\begin{align*}
\text{Bi}_2\text{O}_3 + \text{Nb}_2\text{O}_5 + \text{SrCO}_3 & \rightarrow \text{SrBi}_2\text{Nb}_2\text{O}_9 + \text{CO}_2 \\
\text{BaO} + 6\text{Fe}_2\text{O}_3 & \rightarrow \text{BaFe}_{12}\text{O}_{19}
\end{align*}
\]

Bi₂O₃, Nb₂O₅, SrCO₃ for SBN have been mixed and ball-milled with zirconium balls in a polyethylene container by using a planetary mill for 8 hours, in order to lower it’s grain size and enlarge it’s surface area aiming to achieve more reactive powders. The same calcination conditions were carried out for ferrite for 2 hours. SBN and BFO powders were calcinated at 800°C and at 1000°C for 5 hours respectively.

SBN and BFO powders were mixed as the following rule SBNₓBFO₁₀₀₋ₓ (x = 0, 25, 50, 75 and 100 wt%). Three series of samples were produced for different binders, namely polyvinyl alcohol (PVA), glycerin and tetraethyl orthosilicate (TEOS). The two powders in the defined proportions are mixed with 3 wt% of one of the binders. The powders were used to make cylindrical pellets under an uniaxial pressure of about 0.1 MPa for 5 minutes using a hydraulic press.

The bulks were sinterized at 1050°C for 5 hours. Samples were named in proportion of BFO and SBN (BFO100 -100% BFO, BFO75-75% BFO + 25% SBN, BFO50-50% BFO + 50% SBN, BFO25-25% BFO + 75% SBN and SBN100-100% SBN) and also depending on the binder’s name used during their formation (P, T, G, polyvinyl alcohol, tetraethyl orthosilicate and glycerin respectively). The sintered pellets were polished with fine emery paper in order to make both the surfaces flat and parallel and were electroded with high-purity silver paste for dielectric and electrical measurements to ensure a good electrical contact by using an impedance analyzer to cover the frequency range of interest.

The samples powders were analysed using Mössbauer spectroscopy and X-ray diffraction, both at room temperature. The Rietveld analysis was performed by the refinement program DBWS. The bulk samples were analysed with Raman and infrared spectroscopy and Scanning electron microscope, for investigation of the structure and radiofrequency measurements, as well as, magnetic and electric hysteresis for investigation of the electromagnetic properties.

2.1. X-Ray Diffraction

The X-ray diffractograms were obtained at room temperature, by using non-sinterized powder samples and a Siemens D501 graphite-monochromator diffractometer using a Cu-K radiation detector (k = 0.1542 nm) with a Bragg-Brentano-geometry scintillation counter operating at 40 kV and 25 mA. The ferrite powders were scanned through the 2θ angle ranging from 20° to 85° and the SBN powder ranging from 20° to 70° during five seconds for each step of counting time. The Rietveld analysis was performed by the refinement program DBWS. A Modified Thompson-Cox-Hasting pseudo-Voigt profile function was used to fit the calculated curve into the experimental diffraction data. The obtained density with such refinement was used for calculating the relative density (%) of the samples by the Archimedes method.

2.2. Raman and Infrared Spectroscopy

Infrared spectroscopy was obtained with a Mattson 7000 (FTIR) spectrometer, while Raman spectrum was obtained with a FTIR spectrometer VERTEX 70 equipped with RAM II Bruker FT-Raman module adapted for the rough surfaces (diffusion reflection spectroscopy Easy-
Diff), using bulk sintered samples in both cases.

2.3. SEM

Scanning electron microscope (SEM) images for microstructure observation of the sintered pellets were obtained with a scanning electron microscope Vega XMUT/Tescan, Bruker.

2.4. Radiofrequency Measurements

Both surfaces of pellets were painted with silver paste for impedance at radiofrequency measurements. The permittivity and loss tangent were measured by using Agilent E4992A RF impedance/material analyzer connected to a personal computer, for frequency ranging from 1 MHz to 3 GHz, at room temperature.

2.5. Magnetic Hysteresis Measurements

For the measurement of the magnetic moment (M) versus the applied magnetic field (H) a VSM (vibrating sample magnetometer) was used. All measurements were done at room temperature. Magnetization was normalized by the mass of each sample in order to provide M in emu/g. The samples were measured over a maximum applied field of 8000 Oersted.

2.6. Electric Hysteresis Measurements

The electric hysteresis were obtained from a system composed by a Sawyer-Tower circuit connected to a Agilent 54622A digital oscilloscope, a Agilent 33220A function generator and a Trek 610E high-voltage font. All hysteresis loops were done at room temperature with disc shaped samples.

2.7. Mössbauer Spectroscopy

Mössbauer measurements were performed at room temperature by a FAST (ConTec) Mössbauer Systems spectrometer through transmission geometry and radioactive source of $^{57}$Co in Rh matrix and isometric shifts are related to metallic iron α-Fe. The NORMOS program was used for adjusting iron sites in Barium hexaferrite structure and for determining hyperfine parameters by using powdered samples.

3. Results and Discussion

3.1. X-Ray Diffraction

The X-ray diffraction patterns for BFO100, BFO50 and SBN100 are shown in Figure 1. The diffraction peaks presented by BFO100 sample were identified by JPCDS (file 78-0133), and powder diffraction patterns of BaFe$_{12}$O$_{19}$ located at 33, 16° and 46, 61° were identified by JPCDS (file 20-0132). Some low-intensity peaks observed at 2θ < 30° were attributed to magnetite identified by JPCDS (file 76-0958). Mali and Ataie [4] suggest that the formation of barium hexaferrite comes from the reaction of intermediate phases of iron oxide and a spinel monoferrite BaFe$_2$O$_4$, and the percentage of crystalline BaFe$_{12}$O$_{19}$ grows monotonically as increasing the temperature. This would explain the residual phases and a pure hexaferrite would be obtained for higher calcination temperatures. The presence of residual phases for barium hexaferrite was also reported by other authors [5]. The crystal structure found has hexagonal symmetry belonging to the P63/mmc space group. The average crystallite size was calculated by using the Debye-Scherrer’s equation around 37, 16 nm.

For the SBN100 sample all the diffraction peaks were indexed according to JCPDFS (file 49-0607), without residual phase evidences. The crystal structure found is orthorhombic belonging to the A21am space group with crystalline size estimated about 45 nm.

The diffraction peaks for the present composite shows only a combination of the diffraction of the two base compounds. The representative peaks for SBN100 and BFO100 and spinel monoferrite were indexed by using S, B and F letters respectively.

3.2. Rietveld Refinement

Rietveld Refinement was performed only for the SBN100 and BFO100 samples.
Ferrimagnetism and Ferroelectricity of the Composite Matrix: SrBi₂Nb₂O₉(SBN)x-BaFe₁₂O₁₉(BFO)₁₋₀ₓ

The standard R-factors Rp, Rwp, Rexp, and goodness-of-fit parameter S, as well as the lattice parameters (a, b, c), the unit cell volume (V) and the density are given in Table 1. A low percentage of Rwp and S value close to unit (<1.3 normally) indicates that the refinement was successful. Better results were obtained to SBN100 refinement. The density and volume of unitary cell were close to values reported [3] and also from PDF database.

To the BFO, the presence of a small amount (~0.04%) of secondary phases was not considered in the refinement which increases the refinement parameters values. The density, unit cell and volume presented here are in good agreement with the expected results (JPCDS file peaks number 78-0133) already reported [6,7].

3.3. Archimedes Method

The relative densities of the present samples were measured by Archimedes method and are shown in Table 2. The expected density was achieved and better densification is related to TEOS binder use. The only difference between using TEOS, Glycerin and PVA as binders is the densification factor of the samples, since the organic material should vanish around 500°C and it is supposed not to interfere with the material atomic arrangement. Therefore, no structural differences should be observed, even though electrical and magnetic properties are influenced if any consistent change is observed as causing undesirable effects.

3.4. Raman Spectroscopy

Figure 2 shows the Raman spectra for the composites, SBN100T, and BFOQ25T, BFOQ50T, BFOQ75T and BFOQ100T, respectively. The characteristic bands for SBN were all in agreement with data from the literature [8]. It is usual to all Aurivillius ceramics, once they are peroviskite like materials, a vibration mode over the 800 cm⁻¹ - 900 cm⁻¹ frequency range (associated to an octahedral sub-structure). The 835.95 cm⁻¹ mode was identified to be associated to NbO₆ octahedral stretching mode A₁g. The 585 cm⁻¹ mode is associated to Eₕ modes related to the stretching of oxygen bonds in the octahedron. The lower frequency modes are probably related to lattice vibrations or backscattered photons.

A detailed analysis of hexaferrite vibration modes was carried out by Kreisel [9]. According to his work, 683.7 cm⁻¹ mode was clearly observed, which is related to the symmetric stretching of bi-pyramid shaped sub-structure of FeO₅. This mode is commonly between 670 - 710 cm⁻¹ for M-type ferrite. This is the strongest band in fer-

![Figure 2. Raman spectra for the SBN100T, BFOQ25T, BFOQ50T, BFOQ75T and BFOQ100T samples.](image)

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Ferrimagnetism and Ferroelectricity of the Composite Matrix: SrBi2Nb2O9(SBN)X-BaFe12O19(BFO)100–X

3.5. Infrared Spectroscopy

Figure 3 shows the IR spectra. The NbO6 structure could be related to the 600 - 640 cm⁻¹ bands which are expected for octahedral stretching. The bismuth layers bands are hidden once they would be located about the same frequency range where NbO6 octahedral stretching bands is found, and are wide enough to hide bismuth-layer bands.

Similar results were already reported for the ferrite spectra [10,11-13]. The small displacement of the bands could be related to different methods of sample preparation [11-13]. The identified bands for barium hexaferrite in this present work are 609.4 cm⁻¹, 430 cm⁻¹ and 543.89 cm⁻¹, which are all in good agreement with the characteristic bands for crystalline BaFe12O19 (552, 434 and 583 cm⁻¹) [11-13]. These bands are related to Fe-O vibration bonds. Again, the composites did not apparently show new bands.

3.6. Mössbauer Spectroscopy

The Mössbauer Spectroscopy was carried out only for BFO100, since there is no iron atoms in SBN and the presence of strontium atoms caused great difficulties to obtain the spectra when this material is added into ferrite. Figure 4 shows five observed sextets related to a char-
acteristic complex structure of five iron sites. The sites occupied by iron atoms are 12k, 4f1, 4f2, 2a and 2b, located in the spinel (4f1 and 12k) or hexagonal blocks (4f2, 2a and 2b) which were found in the structure. The hyperfine parameters are given in Table 3 and a fitting method was applied to iron sites based on the same arguments used by Sankaranarayanan and Shaman [14] since our spectrum results show all iron sites formed in the same profile.

The strongest hyperfine field sextet (52, 16T) is associated to 4f2-octahedral site of barium hexaferrite, located on the hexagonal block. The hyperfine fields of 2b (40.65T) and 12k (41.09T) do have close values, being the lowest ones. Although the high quadrupolar splitting associated to 2b makes it possible to identify which one is related to the 2b or 12k site, since 2b site, which has a trigonal bipyramidal symmetry, is highly distorted and is expected to show extremely large quadruple-pole splitting values (2.310 mm/s) and relative intensity (6%) lower than 12k site, which has large relative intensity (>35%) for samples treated at high temperatures (>900°C). Other two sextets have showed high hyperfine fields and they are corresponding to 4f1 (49.00T) tetrahedral and 2a (50.79T) octahedral sites.

Also, the isomer shift is useful for determining valence states, and electron shielding. The oxidation states in which iron form compounds is +3 and +2 and the con-states, and electron shielding. The oxidation states in (50.79T) octahedral sites.

Table 3. Hyperfine parameters of the Mössbauer measurements.

<table>
<thead>
<tr>
<th>Sample</th>
<th>Sites</th>
<th>Coordination</th>
<th>IS (mm/s)</th>
<th>QS (mm/s)</th>
<th>Hhf (T)</th>
<th>RA (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>12k</td>
<td>octahedral</td>
<td>0.351</td>
<td>0.401</td>
<td>41.09</td>
<td>36%</td>
<td></td>
</tr>
<tr>
<td>BFO100</td>
<td>4f1</td>
<td>tetrahedral</td>
<td>0.334</td>
<td>0.097</td>
<td>49.00</td>
<td>18%</td>
</tr>
<tr>
<td>4f2</td>
<td>octahedral</td>
<td>0.372</td>
<td>–0.099</td>
<td>52.16</td>
<td>15%</td>
<td></td>
</tr>
<tr>
<td>2a</td>
<td>octahedral</td>
<td>0.371</td>
<td>–0.098</td>
<td>50.79</td>
<td>24%</td>
<td></td>
</tr>
<tr>
<td>2b</td>
<td>trigonal bipyramidal</td>
<td>0.326</td>
<td>2.310</td>
<td>40.65</td>
<td>6%</td>
<td></td>
</tr>
</tbody>
</table>

Nevertheless most of the grains exhibit multi-domain morphology. BaFe_{12}O_{19} has been reported to have non-isometric crystallite size, which would explain such a wide range of size variation [11]. For the SBN100 the micrographs reveal a much more porous sample, but still dense. The grain size distribution is homogenous and the agglomerate shapes are more symmetric than found in BFO100 samples. The grain size varies between 0.4 - 1 μm, which was also observed by Shrivastava and Jah [2]. The composite sample shows good density with few vacancies. The distribution of ferri-rite and SBN is homogeneous and each grain can be easily identified. The addiction of SBN leads to an increase of the ferrite grain size, which can reach more than 5 μm in composite samples. This change in morphology would probably induce a modification in the magnetic properties.

3.8. Radiofrequency Measurements

As it is shown in Table 4, the values of permittivity (ε’) and loss (tanδ) changes according to different binders. The change in relative density of the sample is related to this effect. The permittivity of BFO samples for different binders is associated to the densification of the sample due to the presence of more vacancies in the lower density samples [17]. The value of the permittivity of this ferrite is slowly decreasing with frequency, because the electronic jump between ferric ions is not able to follow the changes in polarization up to high frequencies [17,18]. SBN is more likely to vary its permittivity as a function of density since the presence of vacancies will induce more significant decreasing in ε’. This is related to the fact that SBN is an insulator, with relative permittivity values going up to 10^2 [2,19].

The differences in density of samples is an important factor for the composite samples. TEOS shows higher densification for the majority of the samples and so higher permittivity values have been observed for BFO75T and BFO50T samples. Figure 6 shows the frequency dependence of permittivity.

Another factor which influences the values of both permittivity and loss is the grain size. Since the addition of SBN is leading to an increase of the ferrite grain size, thus it leads to an easier polarization [20]. The increase of these parameters is observed for composite samples.
Figure 5. SEM micrographs for (a) BFO100; (b) BFO50; (c) SBN100.
with higher percentage of BFO (BFO50 and BFO75).

The dielectric loss shows similar values to TEOS and PVA series and higher values to Glycerin. **Figure 6** shows the frequency dependence of loss tangent (\(\tan\delta\)). The observed behavior for composite samples indicates the existence of resonance absorption at higher frequencies [21]. At low frequency, a weak oscillatory behavior is observed for loss values of SBN. This could be explained by the stronger influence of DC conductivity and charge carriers over this frequency range since no loss peak is observed [22]. The ferrite shows a dispersive behavior for the loss tangent, which decreases as frequency increases. This kind of behavior was studied by Koops [23] and is related to the existence of different regions inside the samples with different values of conductivity. The samples with higher losses were BFO50T and BFO75T.

### 3.9. Magnetic Hysteresis

**Figure 7** shows the hysteresis loops for TEOS and glycerin series. A little depression in the hysteresis loop for pure ferrite samples was observed. This behavior is a consequence of coupling between the barium hexaferrite

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**Table 4. Values of permittivity and loss in the RF range.**

<table>
<thead>
<tr>
<th>Samples</th>
<th>100 MHz</th>
<th>500 MHz</th>
<th>1 GHz</th>
</tr>
</thead>
<tbody>
<tr>
<td>′</td>
<td>tanδ</td>
<td>′</td>
<td>tanδ</td>
</tr>
<tr>
<td>BFO100G</td>
<td>12.38</td>
<td>0.072</td>
<td>12.58</td>
</tr>
<tr>
<td>BFO75G</td>
<td>23.17</td>
<td>0.0374</td>
<td>22.80</td>
</tr>
<tr>
<td>BFO50G</td>
<td>69.75</td>
<td>0.167</td>
<td>48.71</td>
</tr>
<tr>
<td>BFO25G</td>
<td>36.12</td>
<td>0.004</td>
<td>36.84</td>
</tr>
<tr>
<td>SBN100G</td>
<td>35.46</td>
<td>0.048</td>
<td>37.51</td>
</tr>
<tr>
<td>BFO100P</td>
<td>13.19</td>
<td>0.0061</td>
<td>13.43</td>
</tr>
<tr>
<td>BFO75P</td>
<td>77.04</td>
<td>0.0484</td>
<td>78.03</td>
</tr>
<tr>
<td>BFO50P</td>
<td>33.10</td>
<td>0.0854</td>
<td>31.07</td>
</tr>
<tr>
<td>BFO25P</td>
<td>42.42</td>
<td>0.0115</td>
<td>44.02</td>
</tr>
<tr>
<td>SBN100P</td>
<td>27.89</td>
<td>0.0037</td>
<td>28.03</td>
</tr>
<tr>
<td>BFO100T</td>
<td>8.04</td>
<td>0.0055</td>
<td>8.05</td>
</tr>
<tr>
<td>BFO75T</td>
<td>96.68</td>
<td>0.0046</td>
<td>104.8</td>
</tr>
<tr>
<td>BFO50T</td>
<td>108.94</td>
<td>0.128</td>
<td>91.44</td>
</tr>
<tr>
<td>BFO25T</td>
<td>35.35</td>
<td>0.0055</td>
<td>36.47</td>
</tr>
<tr>
<td>SBN100T</td>
<td>39.05</td>
<td>0.0086</td>
<td>39.19</td>
</tr>
</tbody>
</table>

---

**Figure 6.** RF measurements of (a) permittivity and (b) loss tangent for TEOS samples.

**Figure 7.** Magnetic hysteresis loops for (a) TEOS and (b) Glycerin samples.
and the spinel ferrite, which is found for a small percentage and is anti-ferromagnetic. The presence of this extra phase can induce a lowering in the coercivity of the material. This anomalous shape is better observed in the glycerin series loops.

The hysteresis loops show a narrowing behavior by adding SBN. The present material coercivity is reduced as a function of SBN in the composite. And this is coherent because SBN is not magnetic, therefore it is not able to respond to any magnetic stimulation. No response by SBN to magnetic fields causes the lowering in all magnetization parameters. Another important fact is the increase of the ferrite grain size for the composite samples. Depending on the size of the grain, only multi-domains microstructure will be formed and this will lead to a decrease in coercivity and coercive fields, due to the easier domain-wall motion in the multi-domain environment [6,24]. The values of remanence magnetization, saturation magnetization and coercive fields are shown in Table 5. We can see that no strong deviation was observed from TEOS to PVA series. Glycerin series show a little quantitative difference, but follow the same qualitative behavior. Once again, samples density plays an important role of defining the denser samples with higher magnetization parameters and coercive fields. Similar results were reported by Ogasawara and Oliveira [25].

It is known that a good crystalline structure leads to a better magnetization where temperature and dwell time are parameters that strongly influence the enhancement of crystalline structure [6]. Although some works reports high coercivity and coercive fields (>5 KOe) for barium hexaferrite [16,26], this was not observed in this present work because high synthesis temperature used here allows a growth of crystallite size and the formation of multi-domain structures [25]. For the pure ferrite, the higher coercive field observed was 3.7 KOe and the higher magnetization of saturation was 38.3 emu/g, which is in good agreement with similar works [6,18,25].

### 3.10. Electrical Hysteresis

Figure 8 shows P-E hysteresis loops, recorded at room temperature, at 1 Hz frequency, for the selected SBN100P, BFO25P, BFO50P and BFO75P composite samples. The hysteresis loops data are characteristic of a lossy linear capacitor, rather than a ferroelectric material. For the SBN100P the electric field amplitude was gradually increased to show that the polarization saturation tendency is not observed for this sample and as a consequence each hysteresis loop present essentially the same shape, as shown Figure 8(a).

The nonlinear P-E nature of a typical ferroelectric material is characterized by a hysteresis effect, in which the polarization trends to a saturation value for high electric field and the direction of the spontaneous polarization can be reversed by an external applied electric field. On the other hand, a P-E linear behavior is typically observed for normal dielectric or paraelectric materials. As the SBN100P sample was prepared from pressed SBN powders, the P-E behavior observed in Figure 8(a) can be due to typical defects of a non sintered sample. In contrast, the BFO addition at the composite composition is expected to increase the conductivity and consequently a linear P-E behavior could be observed as shown Figure 8(b). Thus, the spontaneous polarization values shown in Figure 8 do not correspond to the absolute values of the studied samples.

### 4. Conclusions

In this paper a study of the magnetic and dielectric properties of composites based on M-type barium hexaferrite BFO (BaFe_{12}O_{19}) and SBN (SrBi_2Nb_2O_9) is presented. The magneto-dielectric matrix composite (SrBi_2Nb_2O_9)(BaFe_{12}O_{19})_{100-x} (x = 0,25,50,75 and 100 wt%) were prepared by a new procedure using the solid state reaction method.

A study of the effects of an organic group of binders

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**Table 5. Magnetic Hysteresis parameters of the samples.**

<table>
<thead>
<tr>
<th>Samples</th>
<th>Remanent Magnetization (emu/g)</th>
<th>Coercive Field (Oe)</th>
<th>Saturation Magnetization (emu/g)</th>
<th>Maximum Field (Oe)</th>
</tr>
</thead>
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Figure 8. Electric hysteresis loops recorded at 1 Hz frequency for (a) SBN100T sample and (b) BFO25P, BFO50P and BFO75P composites.

(TEOS, PVA and glycerin) on the structural properties of the ceramic-matrix composites was also investigated.

SEM micrographics have revealed a good densification and a grain size ordered microns. The relative densities achieved with TEOS as a binder were above 80% for all samples, showing that TEOS is a better binder to this ceramics. It was also observed that density and grain size strongly influences the magnetic and electrical properties, in general, denser samples show better electric and magnetic properties.

A complex behavior was observed for loss tangent over the radiofrequency range, which means that less lossy samples could not help keeping this characteristic over entire frequency range. The magnetic hysteresis loops showed that composite samples preserve the ferrimagnetism for hexaferrite when SBN is added to the composite, although they become less coercive. For electric hysteresis the density of the samples are not high enough to define the true behavior of ferroelectricity in composite samples. For further works, the properties over microwave frequency range, thermal influences on the dielectric properties will be investigated for possible
applications of the specimen.

5. Acknowledgements

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