

Scattering of Conduction Electrons on the Static Deformation Raised by Irradiation in n-GaP Crystals

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Abstract

Study of spectral dependences of absorption coefficient at the region of absorption by free carriers shows that the introduction of radiation defects in n-GaP crystals leads to the appearance of additional scattering besides of traditional ones. A new scattering mechanism on “frozen” phonons (deformation localized near the radiation defects) is suggested and its behavior in experimental and theoretical aspects taking into account X_{1c} - X_{3c} transitions at the scattering process has been studied. It was shown that the increase of “frozen” phonons’ concentration results to the growth of this mechanism contribution in the whole scattering and the absorption coefficient by free carriers is described approximately by law $\alpha \sim \omega^r$, where $-1/2 \leq r \leq 7/2$. Suggested scattering mechanism allows explaining qualitatively the decrease of r . It was established that the disordered by irradiation region effectively decreases the concentration of free carriers, but being a region of increased resistivity, it influences the scattering slightly even at the actual quantum region $\hbar\omega \gg \bar{\epsilon}$.

Keywords

Free Carriers, Absorption, Irradiation, Scattering Mechanism, Deformation Potential, Phonon, Gallium Phosphide

1. Introduction

The basic information about the scattering mechanisms has been obtained from the temperature and field de-

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pendences of kinetic coefficients. In respect of electron-phonon scattering mechanisms, the problem primarily consists in finding the electron-phonon coupling constants in the expression for kinetic coefficients as parameters to be determined experimentally. However, the temperature dependences of kinetic coefficients alone are often insufficient to fully define all of above constants. As shown in [1], even for well-known material such as silicon, it is impossible to consider the scattering mechanisms as finally established. In this regard, the study of frequency dependences of the coefficient of light absorption by free carriers in the quantum frequency range ($\hbar\omega \gg \bar{\varepsilon}$, where $\bar{\varepsilon}$ is the mean kinetic energy of carriers) opens up new possibilities for scientific and applied developments.

Quantum theory of the absorption by free carriers predicts power degree dependences of absorption coefficient:

$$\alpha \sim \omega^{-r}, \quad (1)$$

where r is determined by the scattering mechanisms [2]. Therefore, it can be assumed that different frequency characteristics α are related to different relative contributions of scattering mechanisms. Chemical bond in GaP crystals is of mixed type, so along with the ionized impurities ($r \approx 3 - 35$) both acoustic ($r \cong 1.5$) and optical phonons ($r \cong 2.5$) contribute to the scattering of conduction electrons, and, according to the results of [2], experimentally observed values of degree r should be within the following limits: $3/2 \leq r \leq 7/2$.

For n-GaP usually $r \cong 1.8$ but, as shown in [1] [3], none of the scattering mechanisms, including acoustic-phonon scattering, does not describe absorption of light by free carriers in this crystal. Papers [1] [4] justified the existence of once more type of scattering in n-GaP related to the features of its band structure (so-called i-scattering). However, even though explaining small r for n-GaP, the proposed mechanism does not explain the observed decrease of r with deterioration of the crystal structure perfection, e.g., with increasing compensation ratio [5] [6].

This paper presents the results of study the light absorption by free carriers in n-GaP crystals irradiated by high energy (50 MeV) electrons, *i.e.*, a defects concentration and compensation ratio of free carriers in crystals was carried out by introducing radiation defects. Both experimental and theoretical investigations were applied in the course of electron irradiation at various doses and proposed theory explained the experimental data sufficiently.

2. Experimental Results

The measurements of free carrier absorption in n-GaP crystals doped with tellurium (free carrier concentration $n = (1.5 - 2) \times 10^{17} \text{ cm}^{-3}$) in infrared region ($\lambda = 2 - 15 \text{ }\mu\text{m}$) were carried out using UR-20 spectrophotometer at room temperature. The concentration of charge carriers was determined from the intensity of $3 \text{ }\mu\text{m}$ absorption band based on the value of the absorption cross-section for X_{1c} - X_{3c} transitions, $\sigma = 9 \times 10^{-17} \text{ cm}^2$ [7]. Irradiation by 50 MeV electrons at current density of $1 \text{ }\mu\text{A}/\text{cm}^2 \cdot \text{s}$ at room temperature was carried out using a linear accelerator of Yerevan Physics Institute.

Figure 1 shows the spectral dependence in infrared region of the absorption coefficient of n-GaP crystals irradiated by 50 MeV electrons with various integral fluxes. Absorption in this field is mainly due to: 1) transfer of electrons from the absolute minimum of conduction band X_{1c} to the minimum X_{3c} located higher by 0.4 eV; 2) absorption by free carriers.

The curves reveal an absorption band at $\lambda = 3 \text{ }\mu\text{m}$ (type (1) transition) on the background of uniformly increased with absorption wavelength described by function like (1) due to type (2) transitions. **Table 1** shows the concentration of free electrons, N_A/N_D compensation ratio (N_A -acceptors concentration, N_D -donors concentration), α/n ratio at $\lambda = 6.25 \text{ }\mu\text{m}$ in the initial crystal and after electron irradiation. Also, some data from [5] are presented for identical sample compensated by chemical impurity.

It is seen that as a result of irradiation, the concentration of free electrons decreases and the compensation ratio increases. At the same time, a decrease of degree r in the spectral dependence of the absorption coefficient down to 1.13 takes place. Here, α/n ratio after irradiation increases. When irradiating with 50 MeV electrons, point radiation defects and more thermo-stable disordered regions are simultaneously formed in GaP crystals [8] [9]. **Figure 2** shows the absorption spectra of the crystals: the non-irradiated, irradiated (50 MeV electrons at a dose of $10^{17} \text{ el}/\text{cm}^2$) and after heat treatment at 525°C . For the crystals irradiated with so high integral electron fluxes, the free carriers density decreases sharply, in consequence of which no absorption by free carriers is observed. It was shown in [9] that the point radiation defects in GaP crystals annealed at 525°C , so the differences

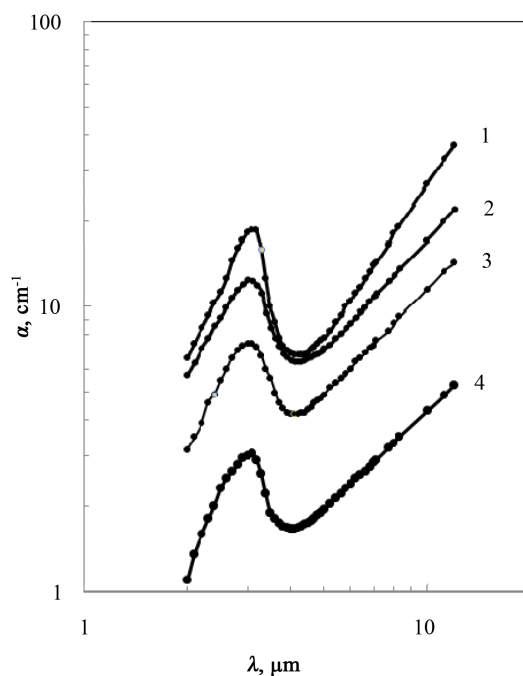


Figure 1. Spectral dependence of the absorption coefficient α for n-GaP sample NO 1 irradiated by 50 MeV electrons with doses: 1- $D = 0$; 2- $D = 4.3 \times 10^{15}$ el/cm²; 3- $D = 7 \times 10^{15}$ el/cm²; 4- $D = 8.6 \times 10^{15}$ el/cm².

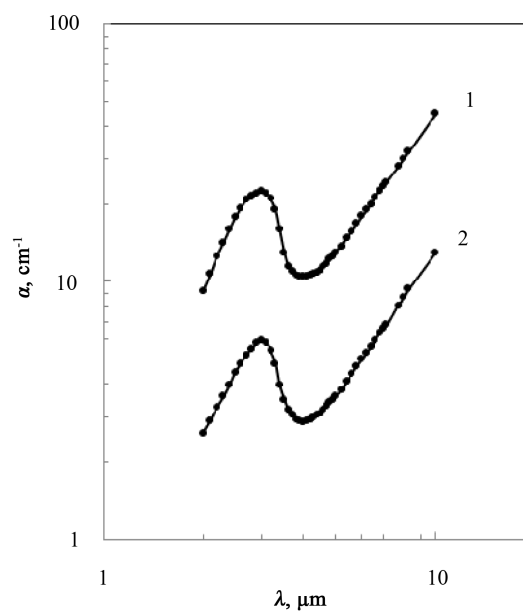


Figure 2. Spectral dependence of the absorption coefficient α for n-GaP sample NO 2: 1-initial; 2-after irradiation ($D = 1 \times 10^{17}$ el/cm²) and annealing at 525°C during 20 min.

between curves 1 and 2 (see **Figure 2**) are due to disordered regions.

It is obvious that whereas only 25% of free electron concentration was recovered after the annealing, the r value in (1) was recovered completely, the α/n ratio also was recovered.

Table 1. The concentration of free electrons.

Numbers of samples	D, el/cm ²	n , cm ⁻³	N_A/N_D	$\alpha/n \times 10^{-17}$, cm ²	r
1	0	1.7×10^{17}	0.1	6	1.75
1	4.3×10^{15}	1×10^{17}	0.4	9.2	1.36
1	7×10^{15}	6.7×10^{16}	0.6	9.7	1.25
1	8.6×10^{15}	2.2×10^{16}	0.87	11.5	1.13
Data from [5]	0	7.2×10^{17}	0.6	20	1.35
2	0	1.9×10^{17}	0.1	10	1.86
2*	1×10^{17}	5×10^{16}	0.7	10.8	1.86

*After irradiation, the sample was annealed at 525°C during 20 min.

3. Theory

The coefficient of absorption by free carriers was calculated by the second order of the perturbation theory [1]:

$$\alpha = \frac{2\pi}{\hbar} \iint \sum_m \frac{|H_{fm}^{cr.i.}|^2 |H_{mi}^{phot}|^2}{(E_m - E_i)^2} f_k (1 - f_{k'}) \delta[E_f - E_i] \frac{2d\mathbf{k}}{(2\pi)^3} \cdot \frac{d\mathbf{k}'}{(2\pi)^3} \quad (2)$$

where i, m, f represent the initial, intermediate and final electron state, respectively, $H_{fm}^{cr.i.}$ is the matrix element of electron interaction with any crystal imperfection ($cr.i.$), H_{mi}^{phot} is the matrix element of the electron-photon interaction, factor $f_k (1 - f_{k'})$ indicates the presence of electrons in the initial and final states. Here, $\alpha(\omega)$ relationship is determined by the dependence of matrix elements on \mathbf{k}, \mathbf{k}' vectors. The expressions $|H_{kk'}^{opt.}|^2, |H_{kk'}^{imp.}|^2, |H_{kk'}^{in.v.}|^2$ describe conduction electron interaction with crystal imperfections [1]. It is shown that $|H_{kk'}^{ac}|^2$ and $|H_{kk'}^{in.v.}|^2$ do not depend on \mathbf{q} ($\mathbf{q} = \mathbf{k}' - \mathbf{k}$), $|H_{kk'}^{opt.}|^2 \sim \frac{1}{q^2}$, $|H_{kk'}^{imp.}|^2 \sim \frac{1}{q^4}$. Substitution of $|H_{mi}^{phot.}|^2$ and $|H_{fm}^{cr.i.}|^2$ to the (2) defines $\alpha(\omega)$ relationship.

In accordance with [1], the frequency dependence of $\alpha(\omega)$ for n-GaP taking into account all traditional scattering mechanisms: on acoustic (ac.) and optical (opt.) phonons, on ionized impurities (imp.), inter-valley (in.v.) scattering; and different types of intermediate states (both intra-band X_{1c} and for overlying X_{3c} zone)-can be represented in the following form:

$$\alpha(\omega) = B \frac{2\beta + 1}{3\beta} \sum_j [A_j^2(\omega) L^j(\beta) + A_{ji}^2(\omega) L^{ji}(\beta)] \quad (3)$$

where $A_j^2(\omega)$ and $L^j(\beta)$ are intensity and anisotropy factor, respectively, for each of the j -scattering mechanisms,

$$\beta = \frac{m_{\parallel}}{m_{\perp}}, \quad B = \frac{4\sqrt{2m_{\parallel}}}{3cn\hbar^2} \frac{e^2 N}{(\hbar\omega)^{3/2}},$$

where N is the carrier concentration, n is the refractive index. The second term in (3) corresponds to the inclusion of intermediate states in the next conduction band (X_{1c} - X_{3c} virtual transitions), which formally looks like the presence of one more scattering mechanism, called i -scattering [1] [4].

The above mentioned experimental results show that the degree r in (1) decreases with increasing N_A/N_D in the case of both doping [5] and irradiation. In both cases the total number of static imperfections increases and around these lattice defects static deformation fields may arise [6]. The additional deformation potential associated with them does not differ, in principle, from the deformation potential caused, for example, by acoustic phonons. The phonons, which correspond to static deformation fields, will be conditionally called “frozen.”

The basic idea of the deformation potential approximation is that the operator of interaction with phonons should be replaced by a quantity characterizing the crystal deformation. It is supposed that the deformation caused by static field arising in the presence of radiation defects, is responsible for local changes in the size and

shape of the crystal unit cell much as in the case of the deformation due to lattice vibrations. If the phonons are the sum of normal vibrations, then the static field near the radiation defects can be considered as the sum of a large number of Fourier components corresponding to the “frozen” lattice vibrations. Obviously, their number is proportional to the density of radiation defects.

Scattering on the “frozen” phonons can be taken into account, on the whole, within the framework of the theory of deformation potential. It is known that in this approximation, when describing the electron interaction with one phonon with wave vector $\mathbf{q} = \mathbf{k}' - \mathbf{k}$, the squared matrix element can be presented in the following form:

$$\left| H_{\mathbf{k}\mathbf{k}'}^{froz.} \right|^2 = \frac{D_q^2 \hbar}{2\rho\omega_{froz.}} \left\{ n_{froz.} + 1 \right. \quad (4)$$

where D_q , $\omega_{froz.}$, $n_{froz.}$ are the interaction constants, frequency and distribution function of such “frozen” phonons, respectively. For small \mathbf{q} ($\mathbf{q} = \mathbf{k}' - \mathbf{k}$) the coupling constant D_q can be expanded to q series:

$$D_q = (D_q)_{q=0} + D_q \left(\frac{\partial D_q}{\partial q} \right)_{q=0} \cdot q \quad (5)$$

and similar to the interaction with the photons, if $(D_q)_{q=0} \neq 0$, the transition is “permitted”, and if $(D_q)_{q=0} = 0$, the electron-phonon transition is “forbidden”.

For the “frozen” phonons, $\omega_{froz.}$ and $n_{froz.}$ weakly depend on q , therefore for the permitted transitions $\left| H_{\mathbf{k}\mathbf{k}'}^{froz.} \right|^2$ does not depend on q , $\alpha_{froz.}(\omega)$ is the same as for the acoustic phonons, so these data cannot explain our experimental results.

However, for the “forbidden” transitions

$$\left| H_{\mathbf{k}\mathbf{k}'}^{froz.} \right|^2 = \frac{\varepsilon_{froz.}^2 \hbar}{2\rho\omega_{froz.}} q^2 \left\{ n_{froz.} + 1 \right. \quad (6)$$

where $\varepsilon_{froz.} \equiv \left(\frac{\partial D_q}{\partial q} \right)_{q=0}$.

For the “forbidden” transitions $(D_q)_{q=0} = 0$, and the electron-phonon interaction will be of the first order in q . Note that for all types of phonons apart from acoustic ones, such scattering does not practically contribute to the kinetic phenomena due small q . However, in the case of the absorption of free carriers, where the scattering is no longer elastic due to presence of high-energy photons $\hbar\omega \gg \kappa_o T$ and $\mathbf{q} \cong \mathbf{k}$, the processes with $\left| H_{\mathbf{k}\mathbf{k}'}^{froz.} \right|^2 \sim q^2$ should be also considered.

Assuming that the interaction with the “frozen” phonons is described by the first order terms in q , with the growth of the number of such phonons the scattering on them can be an essential mechanism competing with the scattering on acoustic phonons.

Below, calculation of $\alpha(\omega)$ for this mechanism is presented similar to the method from [1] for all conventional scattering mechanisms.

Substituting $\left| H_{\mathbf{k}\mathbf{k}'}^{froz.} \right|^2$ and $\left| H_{mi}^{phot.} \right|^2$ in (2) and integrating, we obtain

$$\alpha_{froz.} = \frac{4\pi e^2 N}{nc(\omega)^2} \frac{2\beta+1}{3\beta} \frac{2}{15} \frac{m_{\perp} \sqrt{dm_{\parallel}}}{\pi \hbar^5} \frac{\varepsilon_{froz.}^2 \cdot \chi_{froz.}}{\rho\omega_{froz.}} (\hbar\omega)^{3/2} \left[\frac{2\beta^2 + 11\beta + 2}{2\beta + 1} \right] \quad (7)$$

which can easily be represented in the following form:

$$\alpha_{froz.} = B \frac{2\beta+1}{3\beta} A_{froz.}^2 L_{froz.}, \quad (8)$$

$$\text{where } B = \frac{4e^2 \sqrt{2}}{3ch^2} \frac{N \sqrt{m_{\parallel}}}{n(\hbar\omega)^{3/2}} \quad (9)$$

$$A_{froz.}^2 = \frac{2 m_1 \varepsilon_{froz.}^2}{5 \rho} \left(\frac{\omega}{\omega_{froz.}} \right) \chi_{froz.} \quad (10)$$

$$L_{froz.} = \frac{2\beta^2 + 11\beta + 2}{2\beta + 1} \quad (11)$$

$$\chi_{froz.} = \frac{1}{2} \left[n_{froz.} \left(1 + \frac{\omega_{froz.}}{\omega} \right)^{5/2} + (1 + n_{froz.}) \left(1 - \frac{\omega_{froz.}}{\omega} \right)^{5/2} \right] \quad (12)$$

It follows from the form of $\alpha_{froz.}$ that $\alpha_{froz.} \sim \omega^{-1/2}$ (i.e. $\lambda^{1/2}$), and it is related to the fact that $|H_{kk}^{froz.}|^2 \sim q^2$.

Therefore, the frequency dependence of the total coefficient of absorption by free carriers, taking into account the scattering on the “frozen” phonons, should approximately be described by (1) with $-\frac{1}{2} \leq r \leq \frac{7}{2}$. It should be noted, making allowance for all known scattering mechanisms, that naturally, from a large number of terms not all of them make equal contributions to the final result. In addition, *i*-scattering should be sufficiently intensive as X_{1c} - X_{3c} transition in n-GaP is dipole-allowed. Consideration of all these features allows explaining anomalously low values of *r* in the frequency dependence of $\alpha(\omega)$ and the tendency for a decrease in *r* with increasing the concentration of radiation-induced defects.

4. Discussion and Conclusions

The increase in α/n ratio after irradiation shows that the introduced radiation defects in n-GaP crystals lead to the emergence of an additional scattering mechanism. The change in the scattering mechanism is also indicated by a decrease of *r* in (1) after irradiation.

The comparison of changes in the power degree *r* as a result of introduction of radiation-induced defects with the impact of introduction of compensating impurity atoms on *r* (see **Table 1**) shows that the radiation-induced defects change the scattering mechanism much more efficiently than impurities.

It is known that in $A^{III}B^V$ semiconductor compounds, the radiation defects reveal a significant specific deformation per one point defect which is much more than the specific deformation for a substitution impurity atom. Therefore, the above proposed theoretical conclusion that the scattering on the “frozen” phonons can play a significant role in the absorption by free carriers is explained by additional scattering on local deformations.

In this connection it is interesting to note a much greater influence on the basic absorption edge in GaP, exerted by radiation defects in comparison with impurity centers, which is also explained by deformation effects [8] [9].

The results shown in **Figure 1** reveal a cumulative effect of point radiation defects and disordered regions on the absorption by free carriers in n-GaP.

The relationships shown in **Figure 2** enable to separate the effect of disordered regions on the absorption by free carriers. It follows from **Figure 2** and **Table 1** that the disordered regions both decrease effectively the free carriers concentration and influence on the spectral dependence of the coefficient of absorption by free carriers and on the ratio α/n , and hence influence on the scattering mechanism.

Revealing high potential barrier [9], the disordered regions in n-GaP are non-permeable for majority carriers. Therefore, scattering can only be carried either in peripheral parts of disordered regions, or outside them.

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